

FRB models: (mostly) Robust constraints from observations

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What we would like to know

1. "Progenitors"

- a. The underlying source: NS, solar mass BH, super-massive BH (AGN)...
- b. The energy source: Rotation, B, accretion, planet crashes...
- c. The formation channel

2. Emission mechanism

- a. Dynamics: How is the emitting plasma formed?
- b. Microphysics: Relativistic charged "bunches",
Plasma instabilities ("masers").

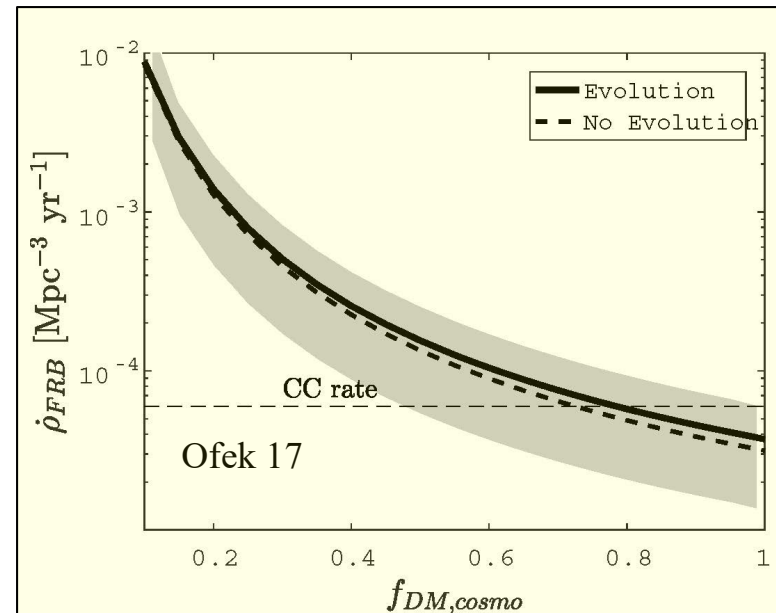
** Progress will be driven by observations, at different rates on different open questions (from 1a to, maybe, 2b).

Duration

- $\Delta t \ll 1 \text{ms}$
 - A natural scale for compact objects
NS, solar mass BH
- [See, however, J. Katz's talk]

Rate

- Uncertainties: small event number, unknown z-evolution & LF.
 - Nevertheless...
 - ~ 3×10^6 /yr @ 1 Jy ms, median DM~800 pc/cm³
 - Assuming IGM dominated DM, DM=250 d_{Gpc}, median d~3 Gpc
- $R \approx \frac{0.5 \times 3 \times 10^6 / \text{yr}}{10^{11} \text{ Mpc}^3} = 2 \times 10^{-5} / (\text{Mpc}^3 \text{ yr})$
- Rate comparable to CC-SN rate,
~100 /yr to 100 Mpc.



- A significant improvement of the constraints may be obtained by understanding the joint ASKAP/Parkes(/CHIME) fluence/DM data.

Energy

- Characteristic radio energy
 $F=1 \text{ Jy ms} = 2 \times 10^{-17} \text{ erg/cm}^2$,
 $E \sim 4\pi(3\text{Gpc})^2 F = 2 \times 10^{40} \text{ erg}$.
(?Parkes/ASKAP: flat LF, Parkes $\sim 10^{-16} \text{ erg/cm}^2$, $E \sim 10^{41} \text{ erg}$).
- X- γ /rays
 $2 \times 10^{-17} \text{ erg/cm}^2 = 10^{-7} \text{ MeV/m}^2$.
LAT [$\sim 30 \text{ MeV}$], GBM [$\sim 1 \text{ MeV}$], BAT [$\sim 30 \text{ keV}$] sub-sec sensitivity
 $\sim 10^{-8} \text{ erg/cm}^2 = 100 \text{ MeV/m}^2 = 10^{46} \text{ erg @ } 100 \text{ Mpc (100/yr)}$.

 \rightarrow Non detection of $\sim 10^2/\text{yr}$ implies $E_{X\gamma} < 10^{46} \text{ erg}$.
Analysis was carried out for LAT [Yamasaki 16], GBM/BAT ?
For the (1st) repeater: $E[\sim 1 \text{ keV}] < 10^{45.5} \text{ erg}$, $E[\sim 30 \text{ keV}] < 10^{47.5} \text{ erg}$.
- Optical
 $10^{46} \text{ erg @ } 3 \text{ Gpc} = 21 \text{ mag [30s]}, 17 \text{ mag [1s]}$
e.g.: ZTF $\sim 21 \text{ mag [1min]}$, $4\pi/10^3$, $0.1\text{T} \rightarrow 100/\text{yr}$.

Energy

- Characteristic radio energy $\sim 2 \times 10^{40}$ erg (10^{41} erg).
- X- γ -O limits
Evidence for $E_{X\gamma} < 10^{46}$ erg
(requires improved LF & z-dist, LAT/GBM/BAT).
Possible optical limit- $E < 10^{46}$ erg.
Improving the limit significantly (/10) will be challenging.
- Radio energy: No strong constraints on the energy source
- Total energy: unknown.
However, $10^{46} (\Delta\Omega/4\pi)$ erg is also not strongly discriminating
(10^{46} erg \sim NS at 1s period).

Duration, Rate, Energy

- $\Delta t \ll 1 \text{ ms} \rightarrow$ likely NS, solar mass BH.
- Rate (1 Jy ms) comparable to CC-SN rate, ~ 100 /yr to 100 Mpc.
- Radio energy $\sim 2 \times 10^{40}$ erg (10^{41} erg),
total probably $< 10^{46}$ ($\Delta\Omega/4\pi$) erg.

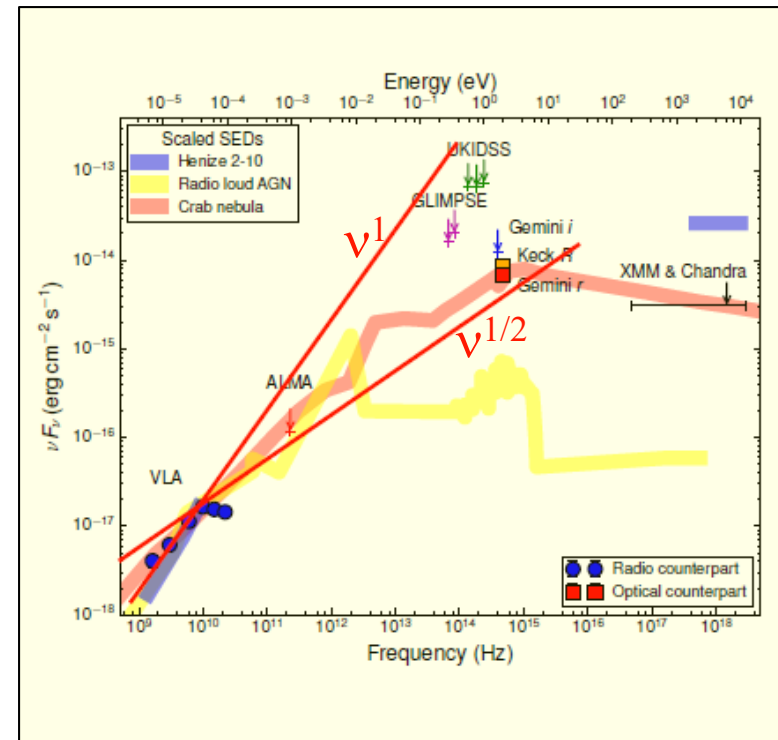
**The persistent radio source associated
with FRB121102:
Robust source constraints**

The persistent radio source associated with FRB121102: Key properties

Assumptions:

- The persistent source was produced in a transient event (no continuous energy output),
- The FRB source is associated with the persistent source, and resides within it.

1. $d_L = 970$ Mpc, $d_A = 680$ Mpc.
2. $t > 4$ yr.
3. $DM = 558 \pm 3$ pc/cm³,
local $DM < 200$ pc/cm³.
4. Angular size consistent with
scatter broadening, $\theta = 0.2 \left(\frac{\nu}{5 \text{ GHz}} \right)^{-2}$ mas.
5. 10 to 30% variability on
10 d time scale at 3 GHz.
6. νf_ν peak 2×10^{-17} erg/cm²s (2×10^{39} erg/s)
at 10 GHz.
7. $\nu f_\nu \sim \nu^1$ down to ~ 1 GHz



Persistent source size

1. Size consistent with scatter broadening
→ $\theta_s \ll 0.2 \text{ mas}$, $R \ll 2 \times 10^{18} \text{ cm}$.

2. ~30% variability at 3GHz on ~10d.

- Intrinsic → $R < 10^{17} \text{ cm}$.

- Refractive scintillation:

$$\theta_d = 0.2 \left(\frac{\nu}{5 \text{ GHz}} \right)^{11/5} \left(\frac{SM_{-3.5}}{80} \right)^{3/5} \text{ mas},$$

$$t_s = 20 \frac{\theta}{0.2 \text{ mas}} \left(\frac{v}{50 \text{ km/s}} \right)^{-1} \frac{d}{1 \text{ kpc}} \text{ day},$$

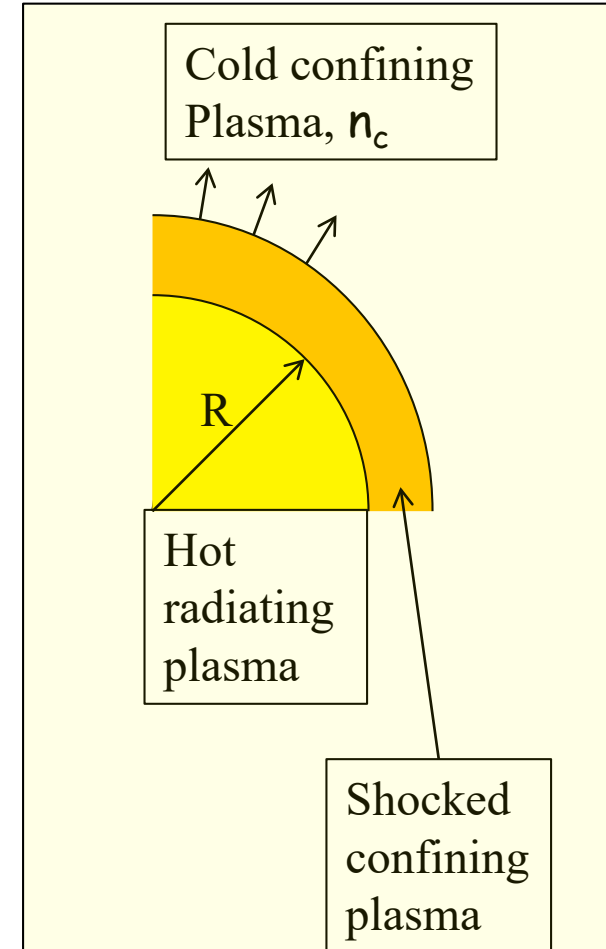
$$\left. \frac{\Delta f}{f} \right|_{\text{max}} = 0.13 \left(\frac{\theta_s}{0.1 \text{ mas}} \right)^{-17/66} \left(\frac{SM_{-3.5}}{80} \right)^{-1/22}$$

$$\text{at } \nu|_{\text{max}} = 3 \left(\frac{\theta_s}{0.1 \text{ mas}} \right)^{-5/11} \left(\frac{SM_{-3.5}}{80} \right)^{3/11} \text{ GHz [e.g. Goodman 97].}$$

→ Variability dominated by scintillation,
 $R < 10^{18} \text{ cm}$.

Persistent source plasma properties

1. $R/t < 10^{10} \text{cm/s}$, no highly relativistic expansion.
 \rightarrow Consider a sphere R of relativistic radiating e^- , with density n_e and magnetic field B .
 2. Peak flux & freq. \rightarrow 2 constraints on $\{R, \gamma_e, n_e, B\}$:
 $B = 10^{-1.5} \gamma_{e,2.5}^{-2} \text{G}$, $n_e = 0.1 \gamma_{e,2.5}^2 R_{17.5}^{-3} \text{cm}^3$
 $(\gamma_e = 10^{2.5} \gamma_{e,2.5})$.
 3. $\nu < 10 \text{GHz}$ spectrum \rightarrow no significant cooling,
 $\rightarrow \gamma_e > 250 (t/10^9 \text{s})^{1/3}$.
 4. No relativistic expansion \rightarrow confining dense plasma, n_c .
 - $v \approx \sqrt{\frac{P}{n_c m_p}} < R/t \rightarrow$ a lower limit to n_c .
 - Shocked shell contribution to DM, δDM ,
 $\nu n_c t, \nu n_c \delta t \propto n_c^{1/2}$, \rightarrow an upper limit to n_c .
- A solution exists- not trivial- for $t < 300 \text{yr}$.



Persistent source plasma properties

1. A solution exists- not trivial- for

$$t < 300 \text{ yr},$$

$$10^{17} \text{ cm} < R < 10^{18} \text{ cm},$$

$$200 < \gamma_e < 10^3, B = 10^{-1.5} \gamma_{e,2.5}^{-2} G,$$

$$E_e \approx 10^{48.5} \gamma_{e,2.5}^3 \text{ erg}, \frac{E_B}{E_e} \approx 1 \gamma_{e,2.5}^{-7} R_{17.5}^3,$$

$$n_c < 10^{2.5} R_{17.5} \text{ cm}^{-3}, M_c < 10^{-1.5} R_{17.5}^4 M_{\text{sun}}.$$

2. Nearly resolved. R may be determined @ 10GHz
(directly if 10^{18} cm , by $\Delta f/f$ if 10^{17} cm).

[Waxman 17]

3. RM & circular polarization?

$$- RM = 1.5 \times 10^5 \text{ rad/m}^2 \rightarrow n_{sc} l_{sc} B_{sc} = 0.2 \text{ pc G/cc.}$$

$$- RM \text{ decrease (30\% @ 1.5yr), DM unchanged,}$$

$$\rightarrow DM_{sc} < 10 \text{ pc/cc} \rightarrow B_{sc} > 0.02 \text{ G.}$$

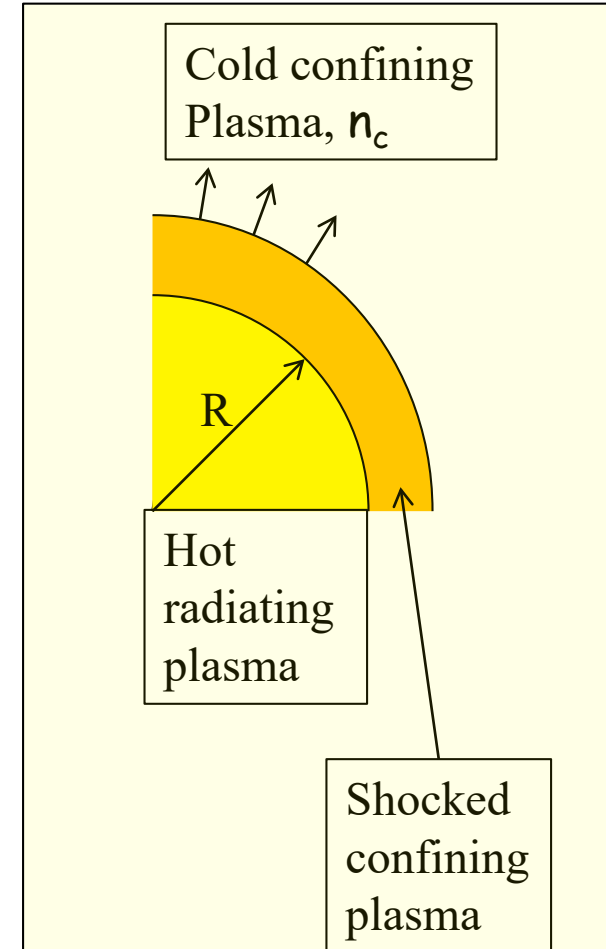
$$- P_{\text{circ}} = RM * \lambda^2 * (\omega_B / \omega) = 0.3 (B_{sc} / 1G) (5\text{GHz} / \nu)^3 < 1\%$$

$$\rightarrow B_{sc} < 0.03 \text{ G} \rightarrow B_{sc} \sim 0.03 \text{ G}, n_{sc} l_{sc} \sim 10 \text{ pc /cc.}$$

** RM screen may be provided by shocked n_c :

$$B_{ep} \sim 0.03 \text{ G}, n_c R \sim 10 \text{ pc/cc.}$$

$$\text{Predicts } P_{\text{circ}} \sim 1 \text{ at } 1 \text{ GHz. [Gruzinov 18]}$$



The persistent source: Progenitor hints

- $\epsilon_e \sim m_p c^2, \frac{E_B}{E_e} \sim 1$ suggests:
Ejection of a mildly relativistic $10^{-5} M_{\text{sun}}$ shell,
that collided with a pre-ejected $M_c \sim 10^{-1.5} M_{\text{sun}}$ shell/"wind".
Possibly: NS formed by accretion induced collapse.
- The plasma properties are similar to those obtained by Beloborodov 17,
Margalit 18. The main differences:
No free parameters and assumptions on age & E_B/E_e ,
which are derived from a Unique over-constrained solution.
- Challenges for meeting the constraints in "Magnetar" models.
 - a. No massive ejecta observed.
 $1 M_{\text{sun}} @ 10^{18} \text{cm}$ produces 30 pc/cc: the bulk of the ejecta must be
driven to large radii, but confining the radiating plasma at 10^{17}cm .
 - b. $N_e \sim 10^{52}$ implies a pair flux (in F_{GJ} units) of $\mu_{\pm} \sim 10^{12}$.
Many orders of magnitude above the Crab, $\mu_{\pm} \sim 10^4$.
Hence, Magnetar models preceding FRB121102 predicted strong
emission peaking at Optical-X rays rather than radio [eg Murase et al. 16].

Producing FRBs

Coherent emission mechanisms

- $E_{\text{FRB}} \sim 10^{39} \text{erg} \sim 10^{-10} E_{\text{persistent}}; \langle L_{\text{FRB}} \rangle \sim 10^{-5} L_{\text{persistent}}$
→ Stringent constraints on the sources are unlikely,
Identification of a unique mechanism is unlikely (see Pulsar history).
- Nevertheless... FRBs are most likely produced by a plasma configuration leading to coherent emission [Katz, Lyutikov...].
- Two types of configurations.
 - a. "Masers": Unstable non-thermal plasma particle distribution, leading to electro-magnetic instabilities.
 - b. "Curvature radiation" from e^- bunches [e.g. Kumar et al 17].

The "gyro-freq." maser

- Reflected electrons form a highly anisotropic momentum distribution, which is unstable.
- Coherent EM waves are produced at the electron gyro-frequency [Sazonov 73; Gallant et al. 92, Lyubarski 14; Ghisellini 17].

- Some difficulties.
 - i. "gyro maser" emission at perpendicular e+ shocks observed in 1D calculations, but suppressed in 2D [Sironi & Spitkovsky 09].
 - ii. In the "magnetar" scenario [Beloborodov 17],

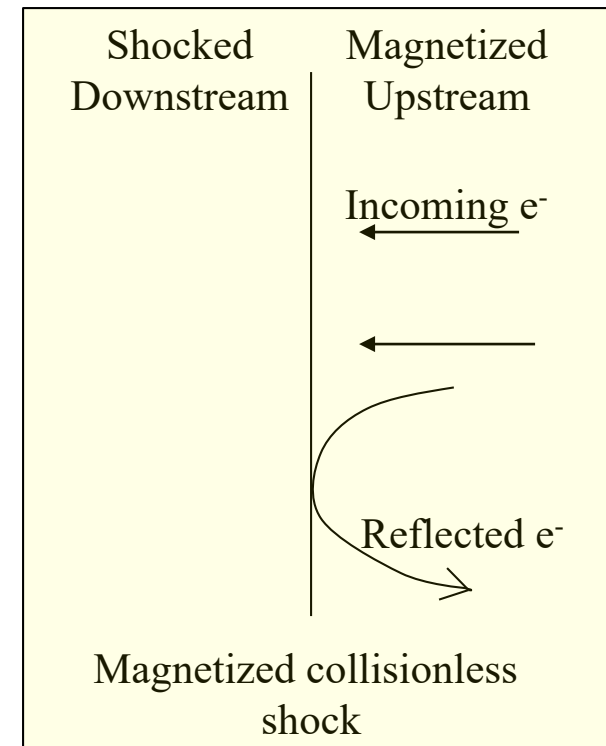
$$\nu = 3 r_{13}^{-1} (L_{f,47} L_{w,39})^{1/4} \text{ GHz},$$

$$E = 10^{40} r_{13} \epsilon_{-1} \sigma_w^{-1} \Gamma_w^{-2} (L_{f,47} L_{w,39})^{1/4} \text{ erg},$$

$$\tau = 3 \Gamma_4^{-2} \mu\text{s}.$$
 7 parameters vs. 3 observed,
 Strong emission at lower radio freq. ($E \sim 1/\nu$).
 - iii. 10^{47} erg at optical/X/ γ :

$$h\nu_{\text{synch.}} = \gamma_e^3 h\nu_{\text{gyro}}$$

$$= 1\text{eV}, 1\text{keV}, 1\text{MeV for } \gamma_e = 10^{2,3,4}$$



Non magnetized collisionless shock maser

- Scattering at the shock leads to isotropic but 'non-thermal' momentum distribution.

If $\frac{dn_e}{d\gamma_e}$ rises faster than γ_e^2 , a maser instability may form [McCray 66; Zheleznyakov 67; Sazonov 70; Sagiv & Waxman 02].

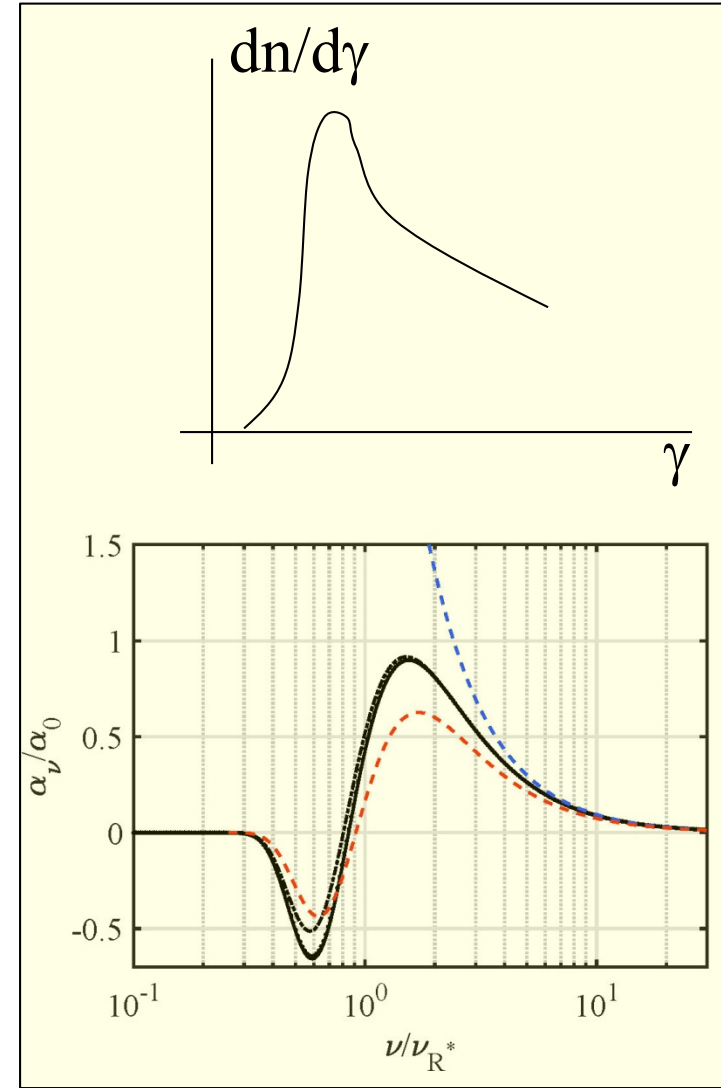
- Instability possible for both $\gamma_e^2 \gg v_p/v_B$

and $\gamma_e^2 \ll \frac{v_p}{v_B} \approx \sqrt{\frac{\epsilon_e}{\epsilon_B}}$ [Sagiv & Waxman 02].

For a narrow e^- distribution and $\gamma_e^2 \gg v_p/v_B$,

$$\alpha_\nu = \alpha_0 F\left(\frac{\nu}{\nu_{R^*}}\right),$$

$$\alpha_0 = \frac{\pi}{2\sqrt{3}} \frac{v_B}{c} \sqrt{\frac{v_B}{v_p}}, \nu_{R^*} = v_p \sqrt{\frac{v_p}{v_B}} \quad [\text{Waxman 17}].$$



Some technical comments

- The maser instability should be derivable directly from a solution of the plasma dispersion relation.

- The Einstein coefficient method provides a "short cut" that is valid for

$$\nu_p^2 \nu_B / \nu^3 \ll \frac{c\alpha_\nu}{\nu} \ll |1 - n| .$$

- These conditions are satisfied at ν_{R^*} as

$$\nu_p^2 \nu_B / \nu^3 \approx (\nu_B / \nu_p)^{5/2} \ll \frac{c\alpha_\nu}{\nu} \approx (\nu_B / \nu_p)^2 \ll |1 - n| \approx (\nu_B / \nu_p) .$$

- A direct solution of the plasma dispersion relation confirms the qualitative results (with some interesting deviations).
- Numerical simulations show that ~1% of the energy may be converted to the maser emission [Gruzinov & Waxman in prep].

Synchrotron maser: Dynamics

- A highly relativistic shell, with energy E_s and $\gamma_s = 10^3 \left(\frac{E_s/10^{41} \text{erg}}{n/0.1 \text{cm}^{-3}} \right)^{1/8} (\Delta t/0.1 \text{ms})^{-3/8}$, is heated by the reverse shock to $T_s \sim m_p c^2$ at $r \sim \gamma_s^2 c \Delta t$ (for source radius $< c \Delta t$).

- In the shocked shell

$$\gamma_s \nu_{R^*} = 0.2 \left(\frac{E_s/10^{41} \text{erg}}{\Delta t/0.1 \text{ms}} \frac{n/0.1 \text{cm}^{-3}}{\epsilon_B/0.01} \right)^{1/4} \text{GHz}$$

$$\alpha_0 \Delta r = 200 \left(\frac{E_s}{10^{41} \text{erg}} \frac{n}{0.1 \text{cm}^{-3}} \frac{\Delta t}{0.1 \text{ms}} \right)^{1/4} \left(\frac{\epsilon_B}{0.01} \right)^{3/4}$$

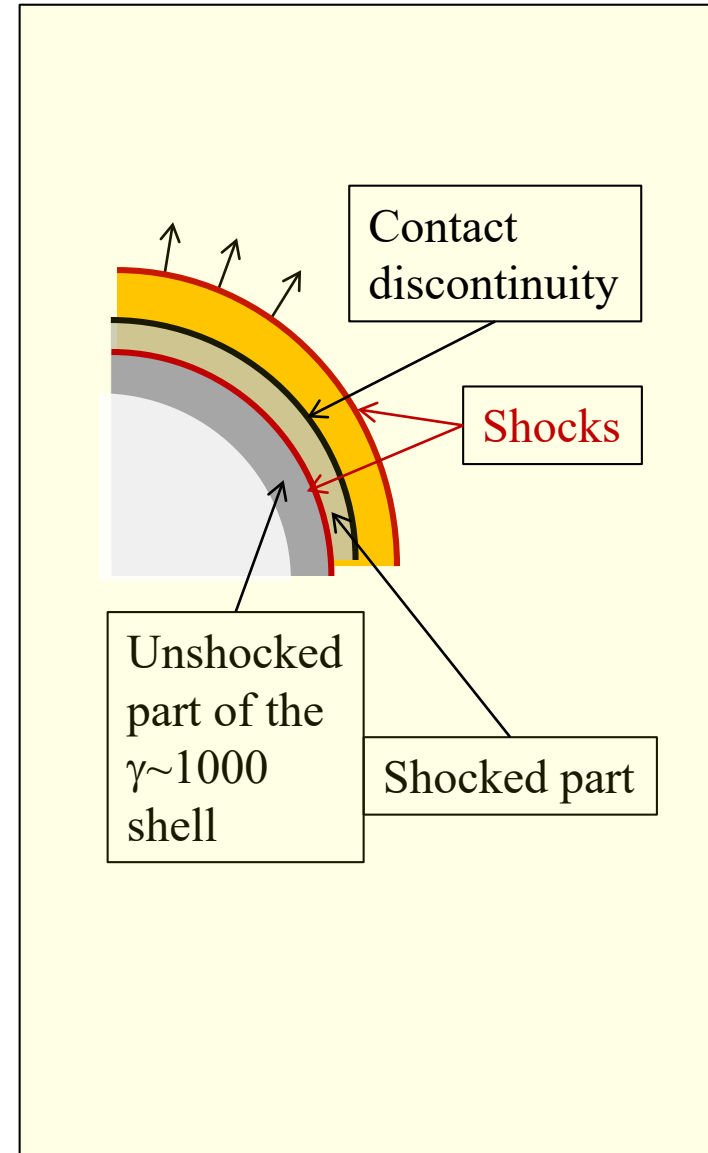
→ For E_s & Δt typical for FRBs:

- $E \sim E_s$ will be emitted over Δt at ~ 1 GHz, provided

$\frac{dn_e}{d\gamma_e}$ is steeper than γ_e^2 below the peak.

- $\{n, \epsilon_B\}$ not free parameters.

- A burst of ~ 10 MeV γ 's with $E < \sim E_{\text{FRB}}$ is predicted from the forward shock.



Summary

- $\Delta t \sim 1\text{ms} \rightarrow$ likely compact- NS, solar mass BH.
- Rate (1 Jy ms) comparable to CC-SN rate, ~ 100 /yr to 100 Mpc.
All probably repeating (0.1 Jy ms rate).
- Radio energy $\sim 2 \times 10^{40}$ erg (10^{41} erg), total probably $< 10^{46}$ ($\Delta\Omega/4\pi$) erg.
- ** Rate & E constraints may be improved by
ASKAP/Parkes/Chime LF & z-dist. and Opt(-X- γ) observations.
- ** Progenitors will most likely be identified by localization.

- (1st) Persistent source:
 $t < 300$ yr, $E = 10^{49.5}$ erg in $10^{-5} M_{\text{Sun}}$ surrounded by $< 10^{-1.5} M_{\text{Sun}}$ at $\sim 10^{17}$ cm.
Hints to a NS formed with relatively low M & E ejecta, AIC.

- FRB mechanism- most likely plasma instability ("maser"), identification of a unique instability is unlikely (see Pulsar history).
 - Gyro-maser: $E \sim 1/\nu$ problem, predicts $E \sim 10^{46}$ erg in optical/X/ γ .
 - Non-magnetized maser: efficient conversion of kinetic energy to coherent radio emission, weak ~ 10 MeV emission.