

The price of tuning

Unique probes of tuned axions

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Stanford

Tuning

- Usually Tuning is a bad thing
 - Naturalness has guided particle physics thinking for a very long time
- What if we're wrong?
 - Is this a bad thing?
- No!
- Interesting signatures can result from tuning

Axion solution

$$\mathcal{L} \supset \frac{g^2}{32\pi^2} \left(\theta - \frac{a}{f_a} \right) G_{\mu\nu} \tilde{G}^{\mu\nu} + \frac{1}{2} \partial_\mu a \partial^\mu a$$

- One parameter solution (KSVZ axion)
- Also a dark matter candidate
- String theory motivation for not just one axion, but many many axions

Axion solution

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Axion dynamically sets
the neutron EDM to 0

$$V = -m_\pi^2 f_\pi^2 \left| \cos \left(\frac{a}{2f_a} \right) \right|$$

Axion parameter space

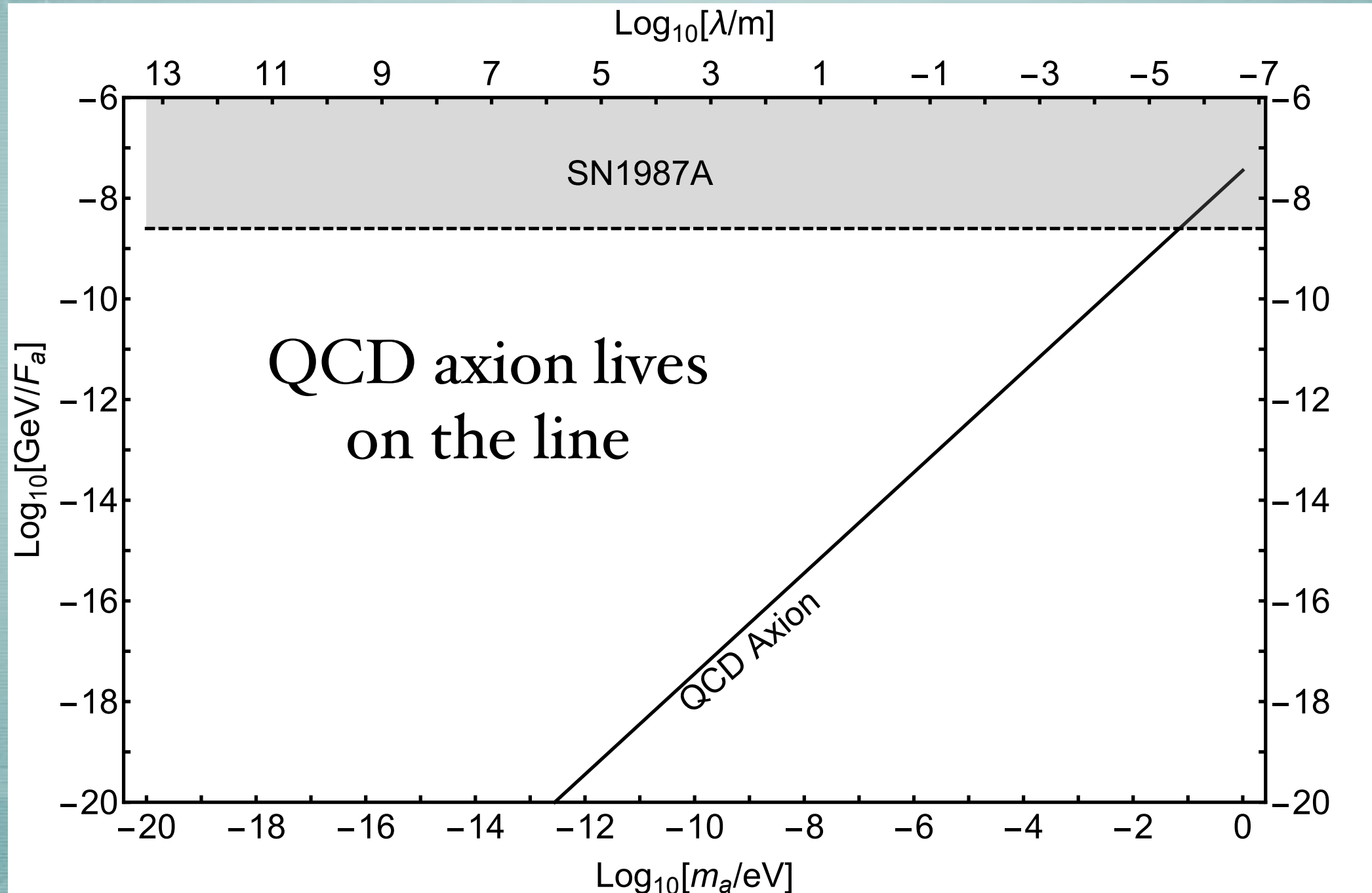
- Instead of

$$V = -m_{\pi}^2 f_{\pi}^2 \left| \cos \left(\frac{a}{2f_a} \right) \right|$$

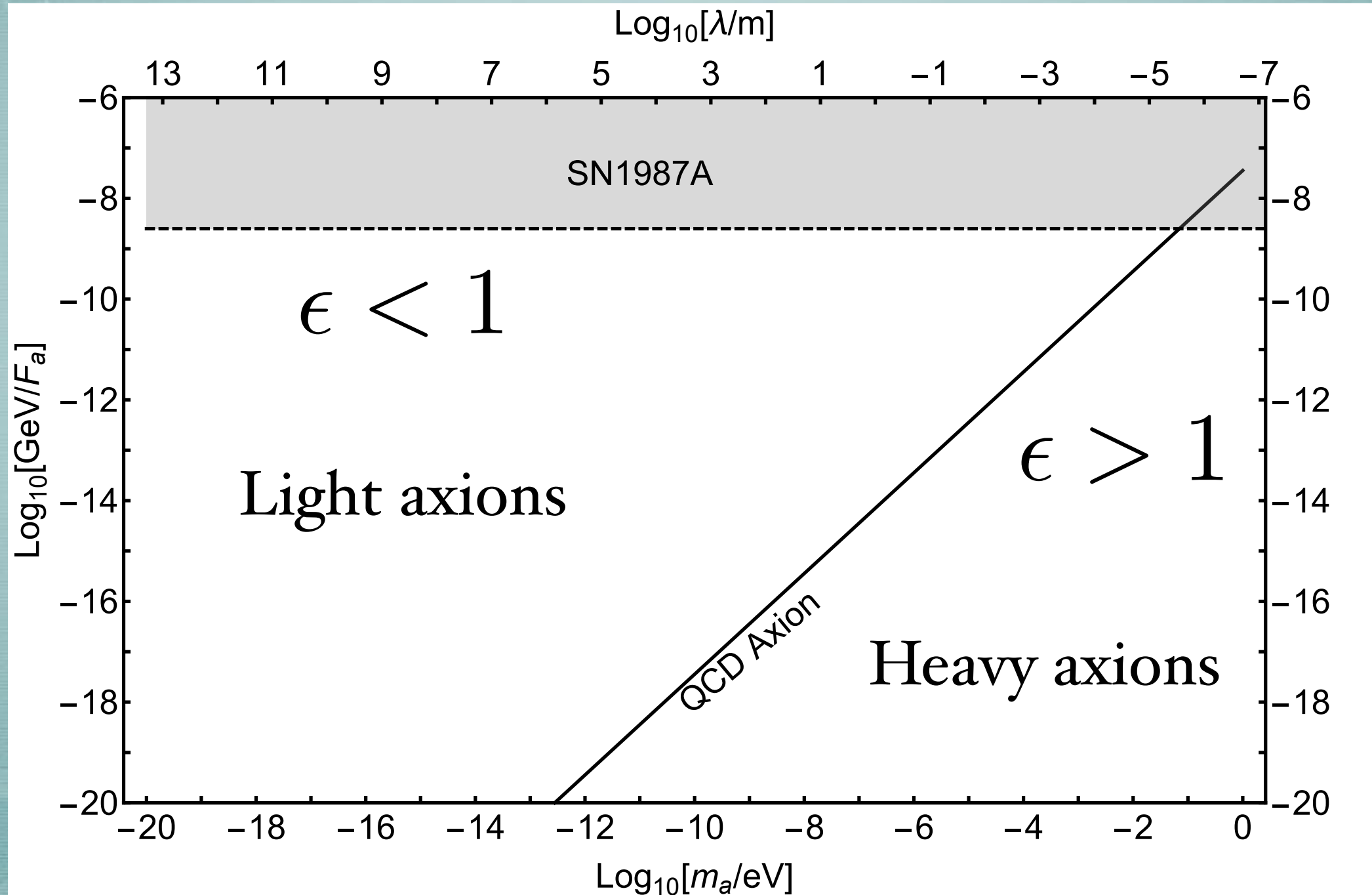
- We have

$$V = -\epsilon m_{\pi}^2 f_{\pi}^2 \left| \cos \left(\frac{a}{2f_a} \right) \right|$$

Axion parameter space



Axion parameter space



Tuning

$$\mathcal{L} \supset y\phi\bar{\psi}\psi$$

- The biggest difference between tuned particles and non-tuned particles is sensitivity to finite density corrections

$$m_\phi \sim y\Lambda$$

$$m_\phi \sim yT$$

- If tuned, thermal mass can be larger than bare mass

Neutron Stars + axions

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- Large finite density objects
- How does the axion behave around and in large density objects?

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- Large finite density objects
- How does the axion behave around and in large density objects?
- At high density, QCD deconfines so finite density makes the QCD contribution to the axion potential go away

Neutron Stars + axions

- Axion potential depends on quark condensate

$$m_u(\langle \bar{u}u \rangle_{n_N} - \langle \bar{u}u \rangle_0) = -m_u(\langle \frac{\partial H}{\partial m_u} \rangle_{n_N} - \langle \frac{\partial H}{\partial m_u} \rangle_0)$$

Neutron Stars + axions

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$$H = E = m_n n_N$$

Neutron Stars + axions

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$$\begin{aligned} m_u (\langle \bar{u}u \rangle_{n_N} - \langle \bar{u}u \rangle_0) &= -m_u \left(\left\langle \frac{\partial H}{\partial m_u} \right\rangle_{n_N} - \left\langle \frac{\partial H}{\partial m_u} \right\rangle_0 \right) \\ &= -m_u \left\langle \frac{\partial E}{\partial m_u} \right\rangle = \sum_{N=n,p} n_N \sigma_N^u \end{aligned}$$

$$H = E = m_n n_N \qquad \sigma_N^u \equiv m_u \frac{\partial m_N}{\partial m_u}$$

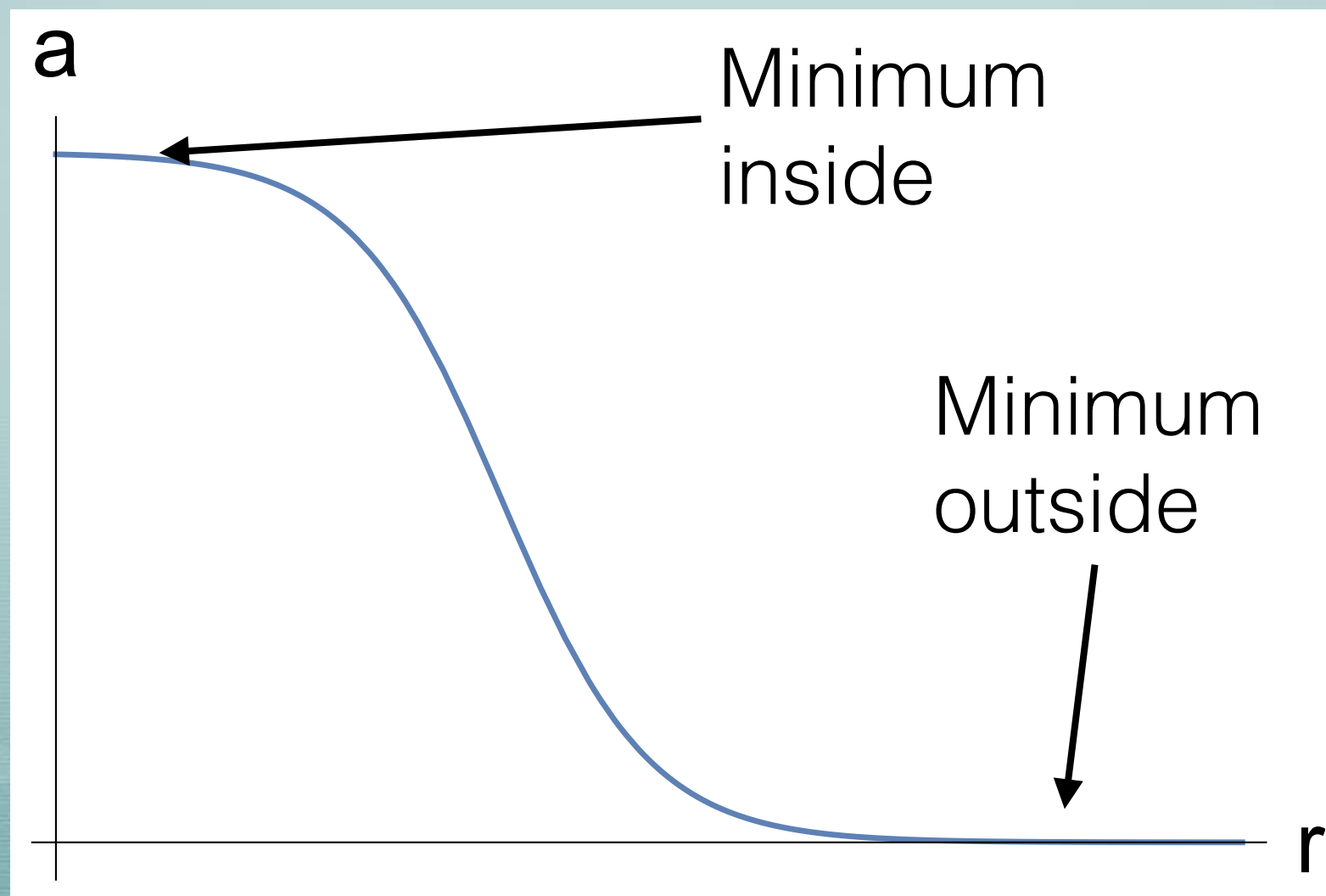
Neutron Stars + axions

$$V = -m_\pi^2 f_\pi^2 \left(\epsilon - \frac{\sigma_N n_N}{m_\pi^2 f_\pi^2} \right) \left| \cos \left(\frac{a}{2f_a} \right) \right| + \mathcal{O} \left(\left(\frac{\sigma_N n_N}{m_\pi^2 f_\pi^2} \right)^2 \right)$$

- If object is dense enough, sign of the potential flips when perturbation theory still valid!
 - Part of the reason why we needed a slightly tuned axion
- Neutron stars can source the axion!

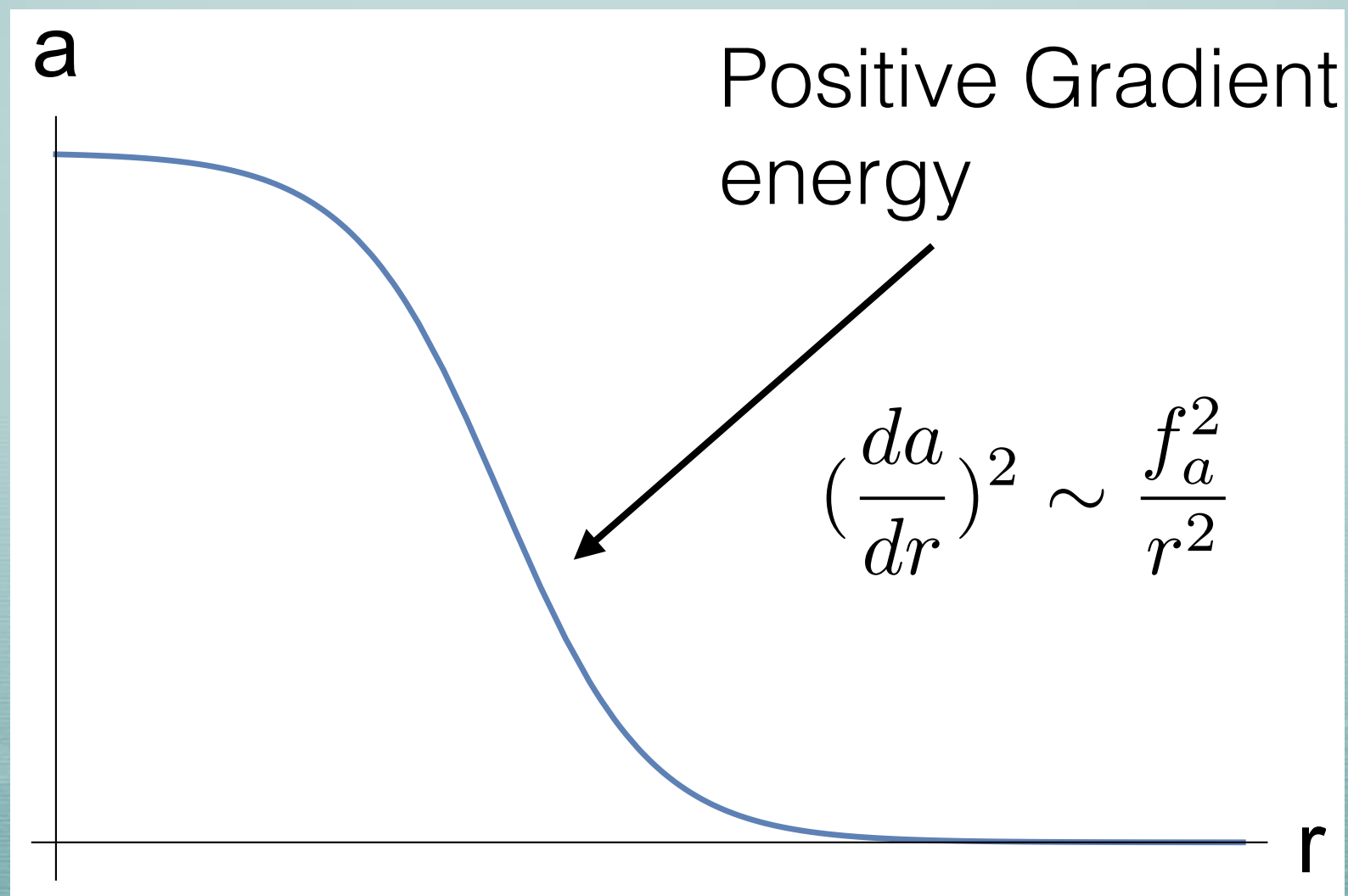
Neutron Stars as axion sources

- When does this happen?



Neutron Stars as axion sources

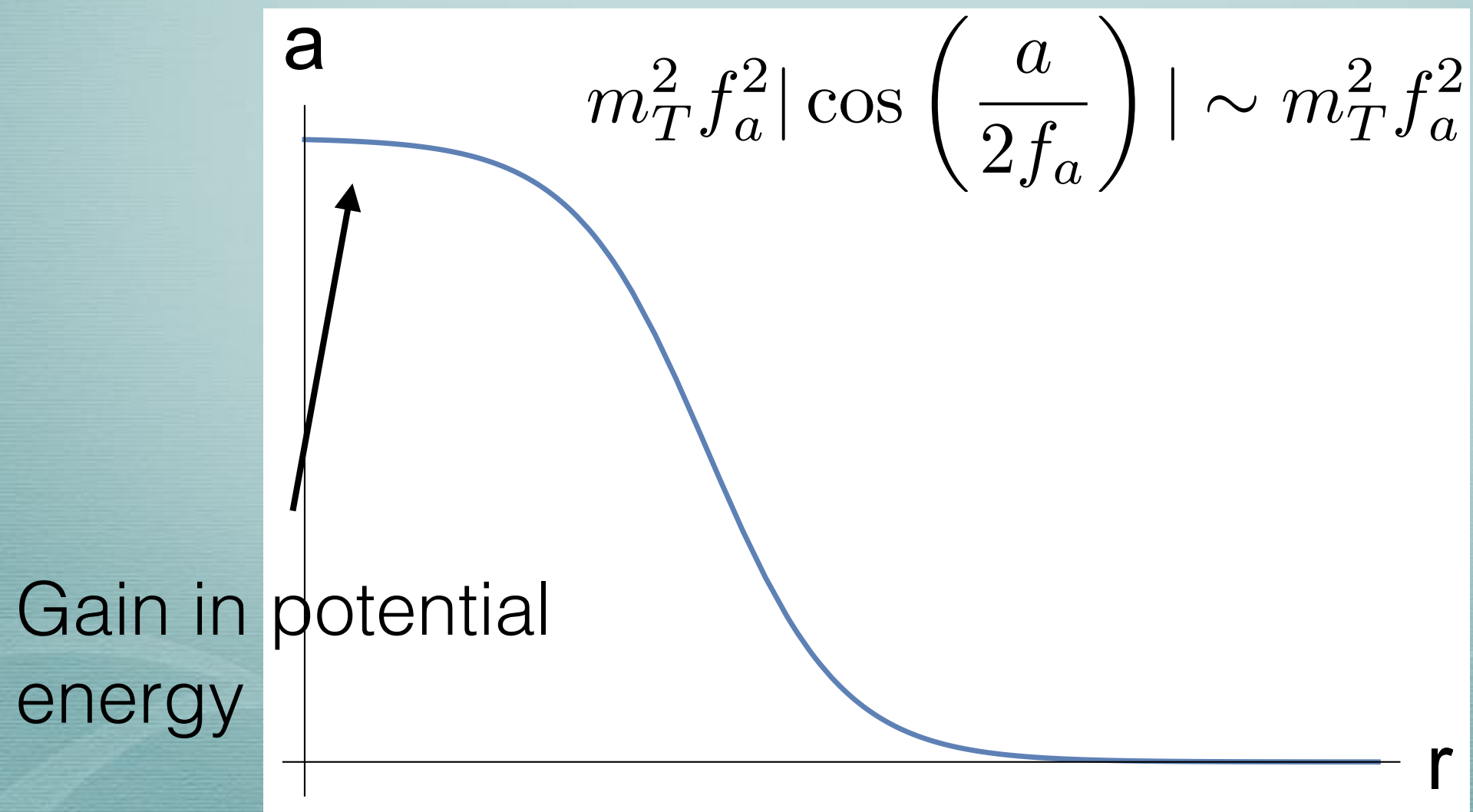
- Energy!



Neutron Stars as axion sources

- Energy!

$$m_T = m_\pi f_\pi \frac{\sqrt{\frac{\sigma_N n_N}{m_\pi^2 f_\pi^2} - \epsilon}}{2f_a}$$



Neutron Stars as axion sources

$$\Delta E \sim \text{Volume} \times \left(\frac{f_a^2}{r^2} - m_T^2 f_a^2 \right) < 0$$

- Critical radius at which neutron stars start to source the axion

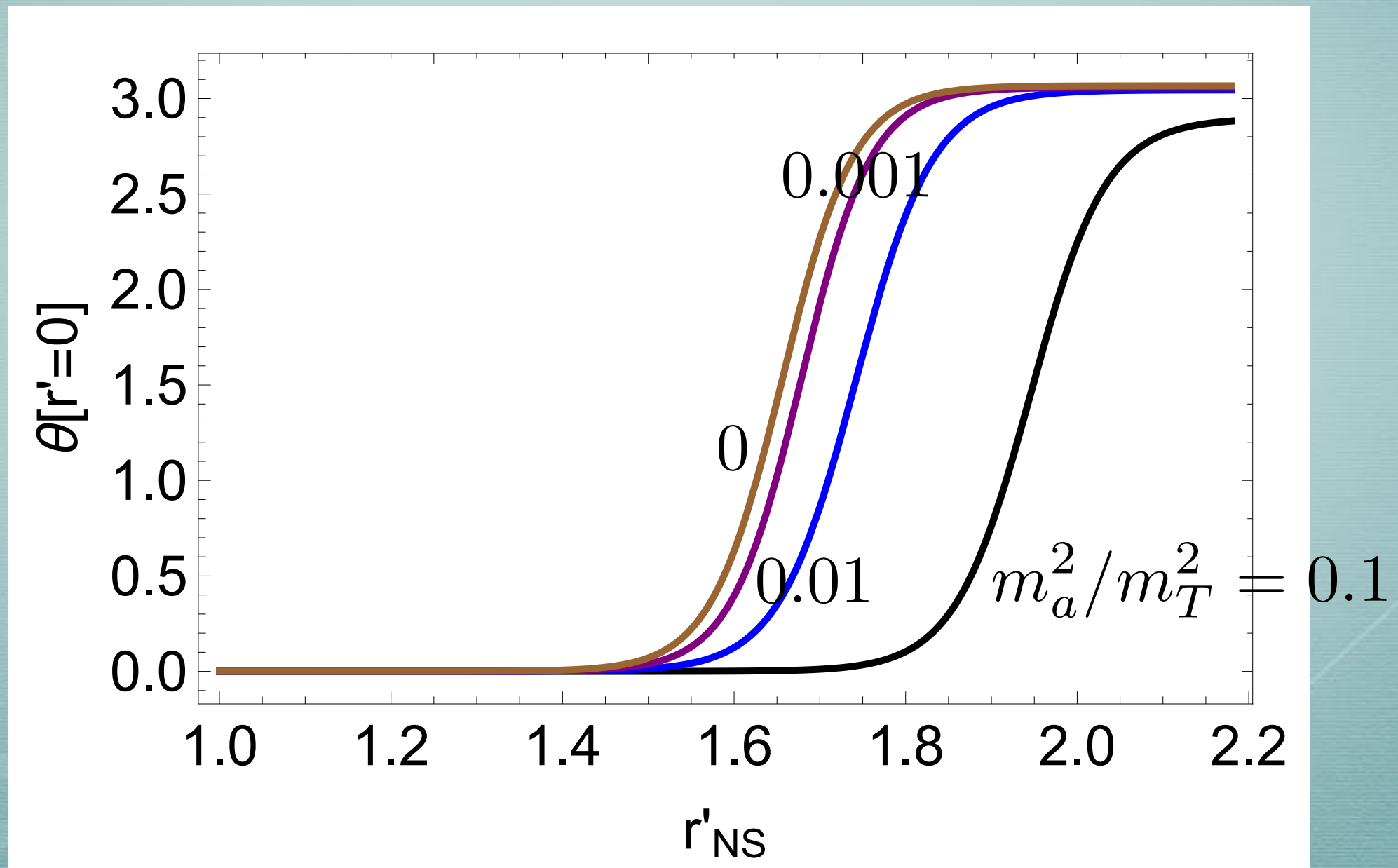
$$r_{\text{crit}} \gtrsim \frac{1}{m_T}$$

Neutron Stars as axion sources

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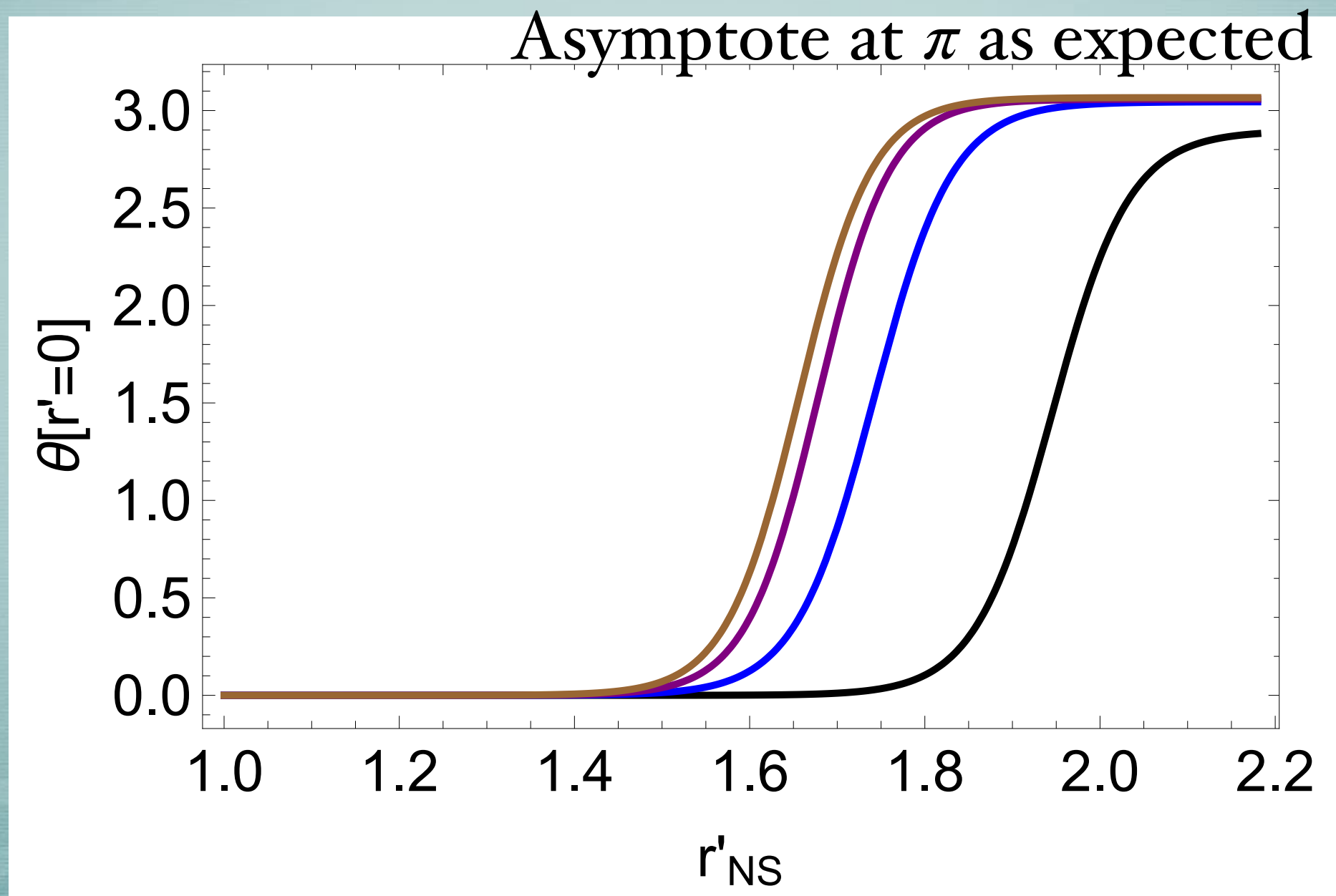
- Both neutron stars and nuclei are nuclear densities
- Only neutron stars large enough to source the axion
- Nuclei too small!
 - Fifth force constraints do not apply

Neutron Stars as axion sources



$$r' = r m_T$$

Neutron Stars as axion sources



$$r' = r m_T$$

Axionic Force

- What happens when two objects that source the axion get close to each other? (Hint : the electron sources the electromagnetic field)

Axionic Force

- What happens when two objects that source the axion get close to each other? (Hint : the electron sources the electromagnetic field)
- There is a force!
- The axion mediates a force between neutron stars!

Axionic Force

$$a = q \frac{e^{-m_a r}}{4\pi r}$$

- Because the axion has a field value about π inside of the neutron star

$$q \sim 4\pi^2 f_a r_{NS}$$

- At long distances, easy to show that the force is a standard Yukawa force

$$V = -q_1 q_2 e^{-m_a D} / 4\pi D$$

Axionic Force

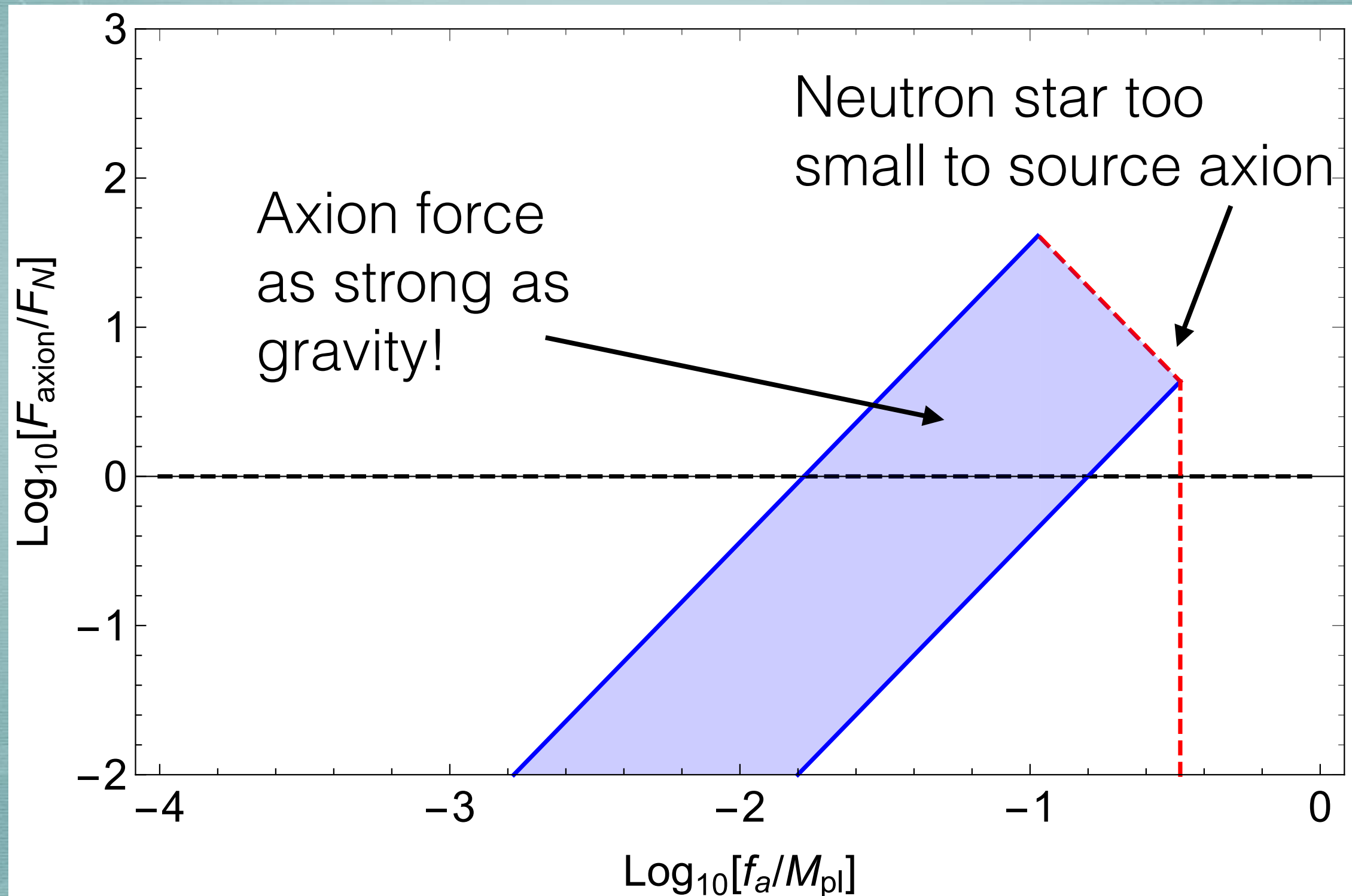
$$V = -q_1 q_2 e^{-m_a D} / 4\pi D$$

- Attractive or Repulsive interaction

$$V = -m_\pi^2 f_\pi^2 \left(\epsilon - \frac{\sigma_N n_N}{m_\pi^2 f_\pi^2} \right) \left| \cos \left(\frac{a}{2f_a} \right) \right| + \mathcal{O} \left(\left(\frac{\sigma_N n_N}{m_\pi^2 f_\pi^2} \right)^2 \right)$$

- Due to symmetry, for every solution (charge) there is a negative solution

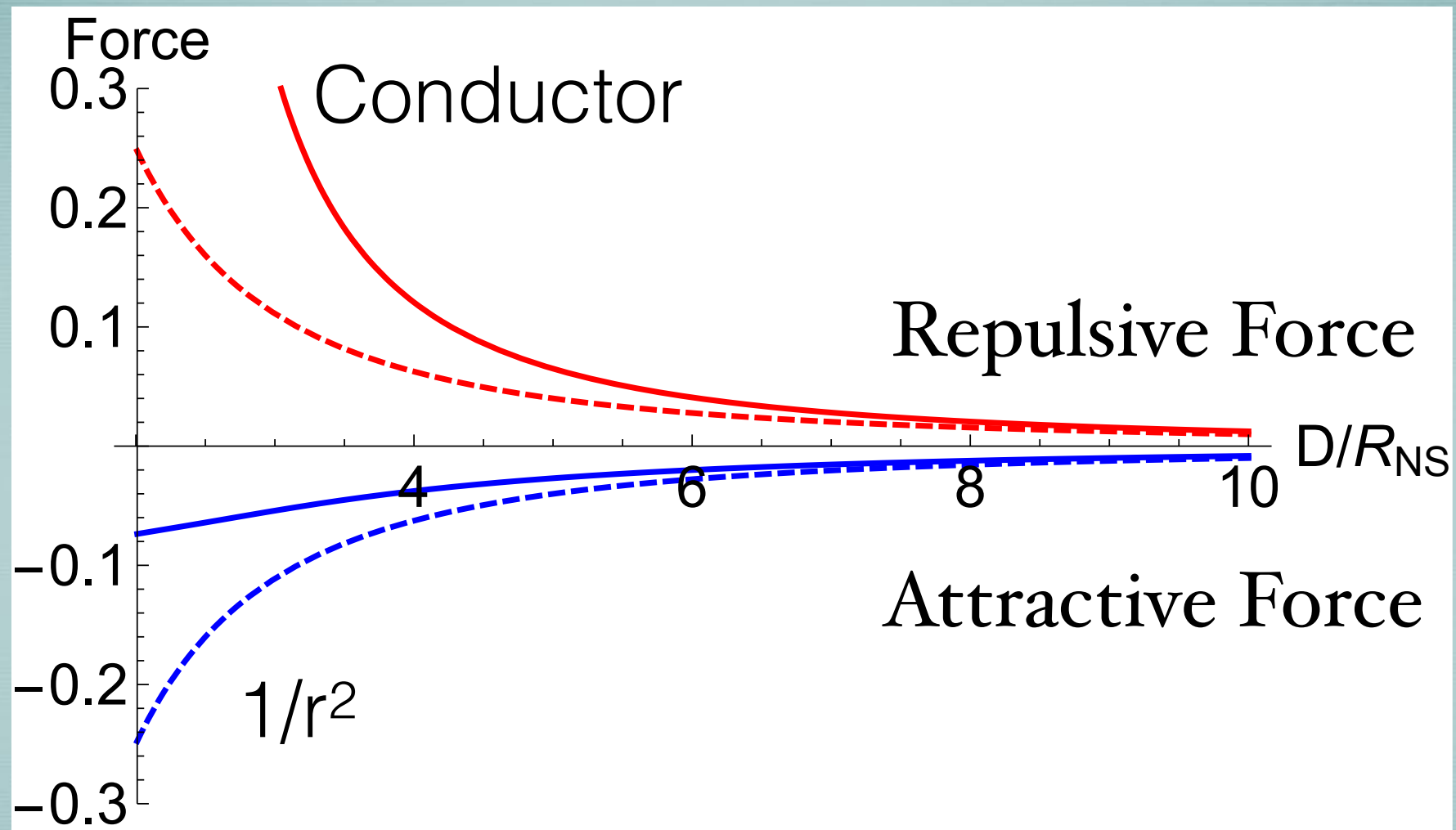
Axionic Force



Axionic Force

- Axion force is actually a boundary condition problem
 - In the limit where the tachyonic mass is infinite
- The axion is stuck to a particular value at the surface of the neutron star
- This is just like a conductor!
- The force between neutron stars can be calculated just like in E+M with a method of images

Axionic Force



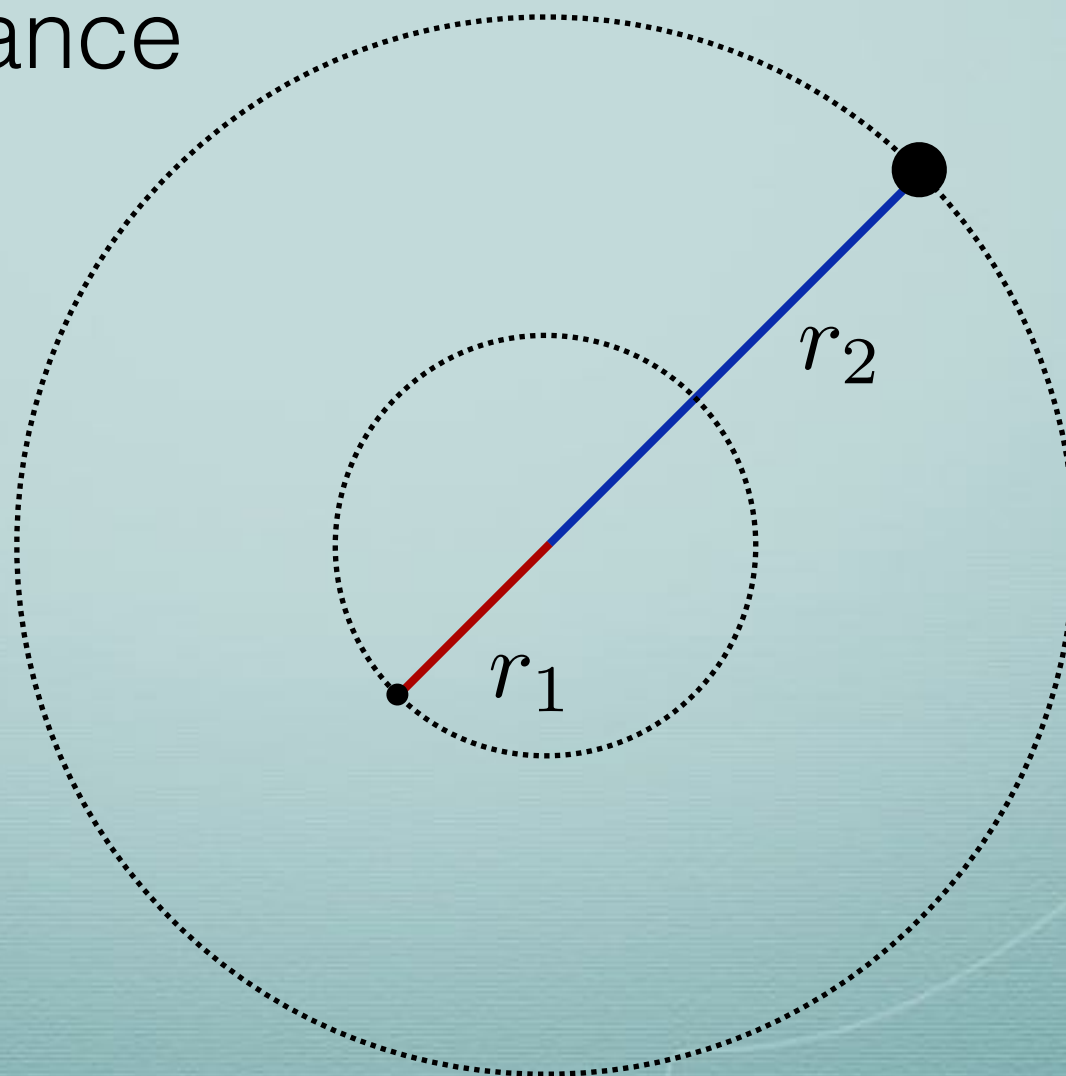
- Not a standard $1/r^2$ force at short distances
- Attractive force weaker
- Repulsive force stronger

LIGO implications

- LIGO has seen merging neutron stars
- What sort of implications will this have for neutron stars inspirals?
 - Mergers will need numerical simulations
- New force
 - At $r \sim 1/m$, the new force kicks in and increases or decreases the frequency
- New radiation
 - When frequency $\sim m$, new scalar Larmor radiation

Inspiral

- Inspiral can be calculated in a simple analytic way
- Assume instantaneously circular orbit with time dependent distance



Inspiral

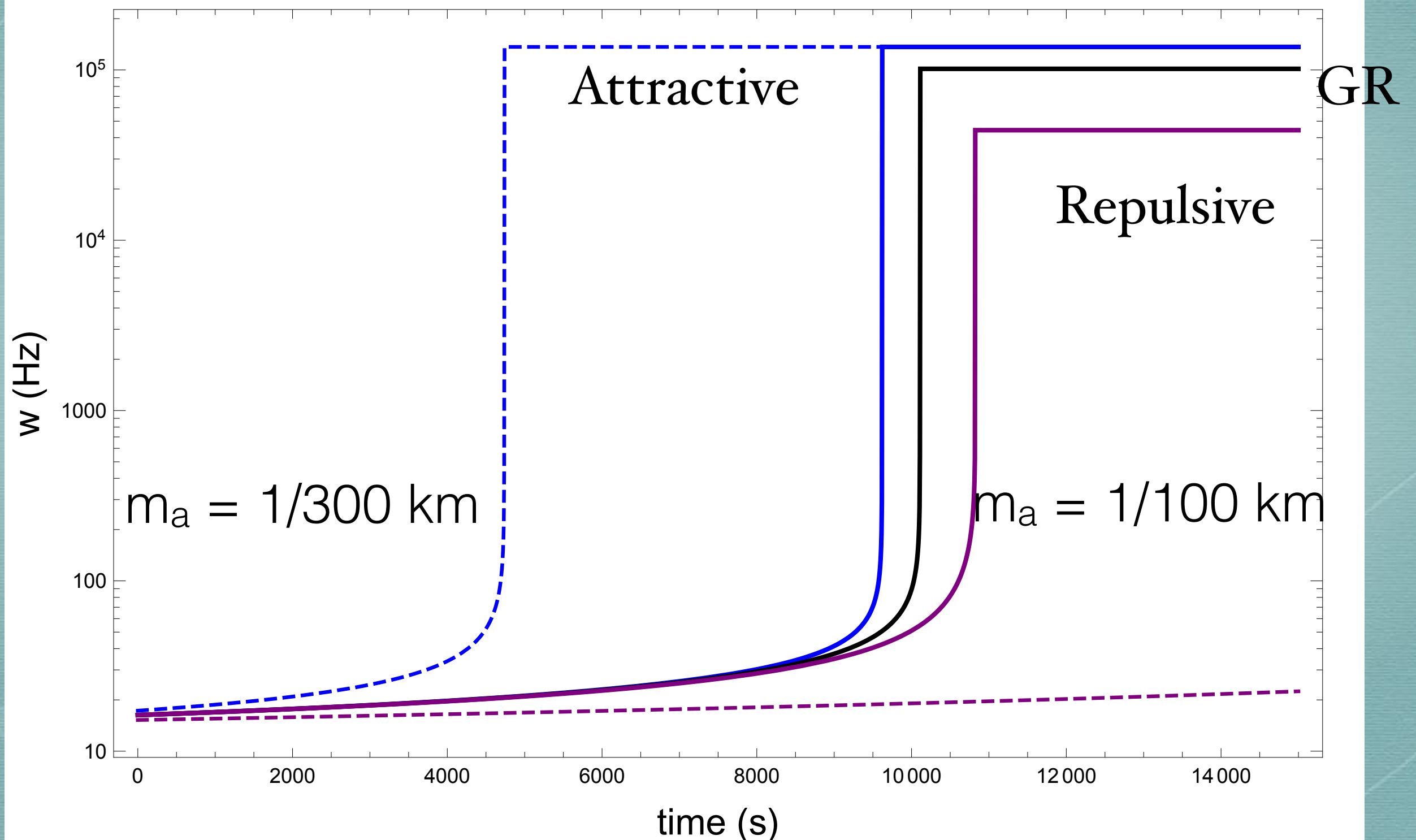
- Power emitted comes from gravitational quadrupole radiation and scalar dipole radiation

$$\frac{dE}{dt} = -\frac{32}{5}G\mu^2 D^4 \omega^6 - \frac{1}{4} \frac{\omega^4 p^2}{6\pi} \left(1 - \frac{m_a^2}{\omega^2}\right)^{3/2} \Theta(\omega^2 - m_a^2)$$

$$p = q_1 r_1 - q_2 r_2$$

- Can calculate everything we want about the early inspiral phase from here

Inspiral - 80% gravity



Inspiral

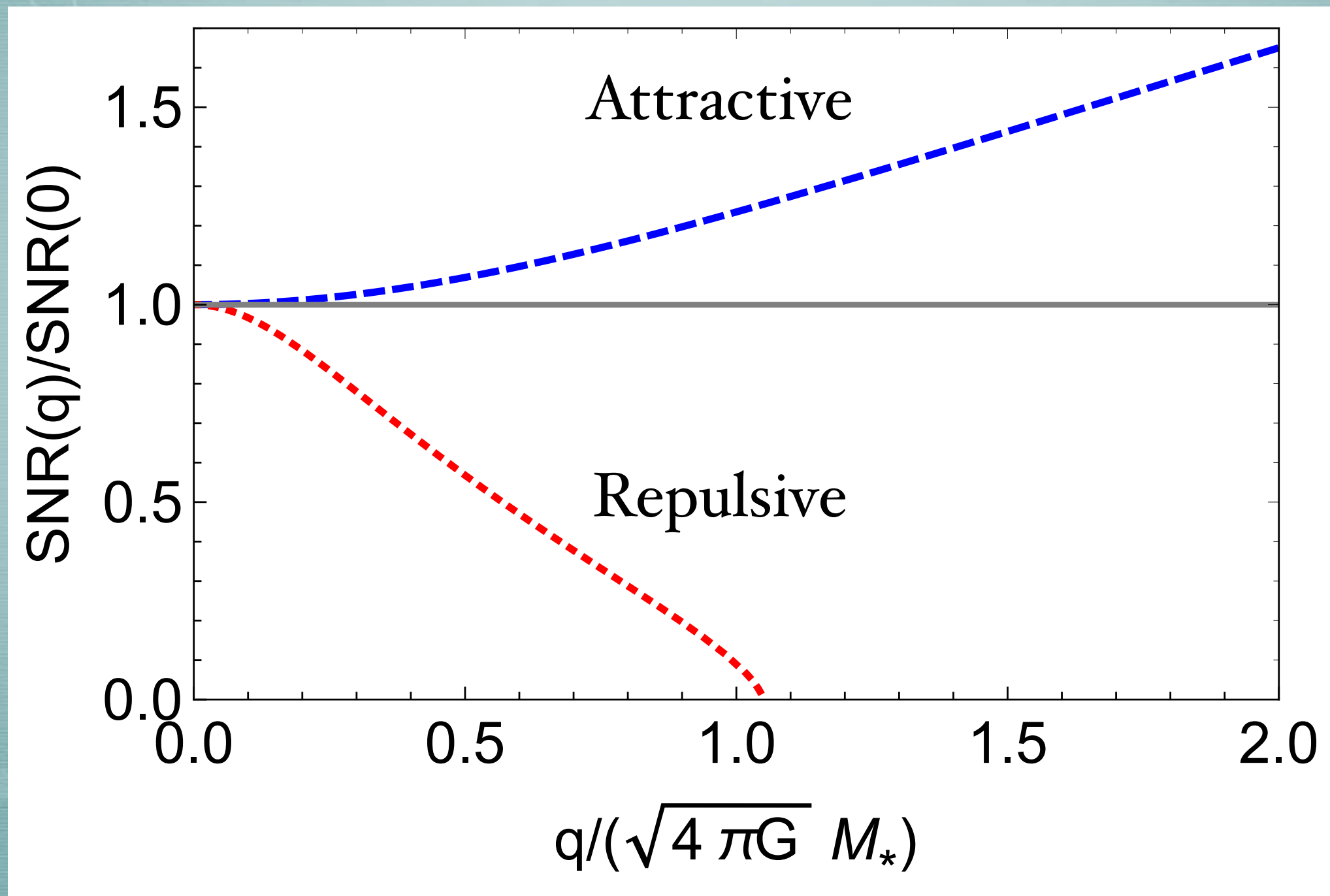
- Two major effects
- Quality factor
 - Increases : Repulsive Interactions
 - Decreases : Attractive Interactions
- Amplitude
 - Decreases : Repulsive Interactions
 - Decreases : Attractive Interactions

$$Q \sim \frac{\omega^2}{\dot{\omega}}$$

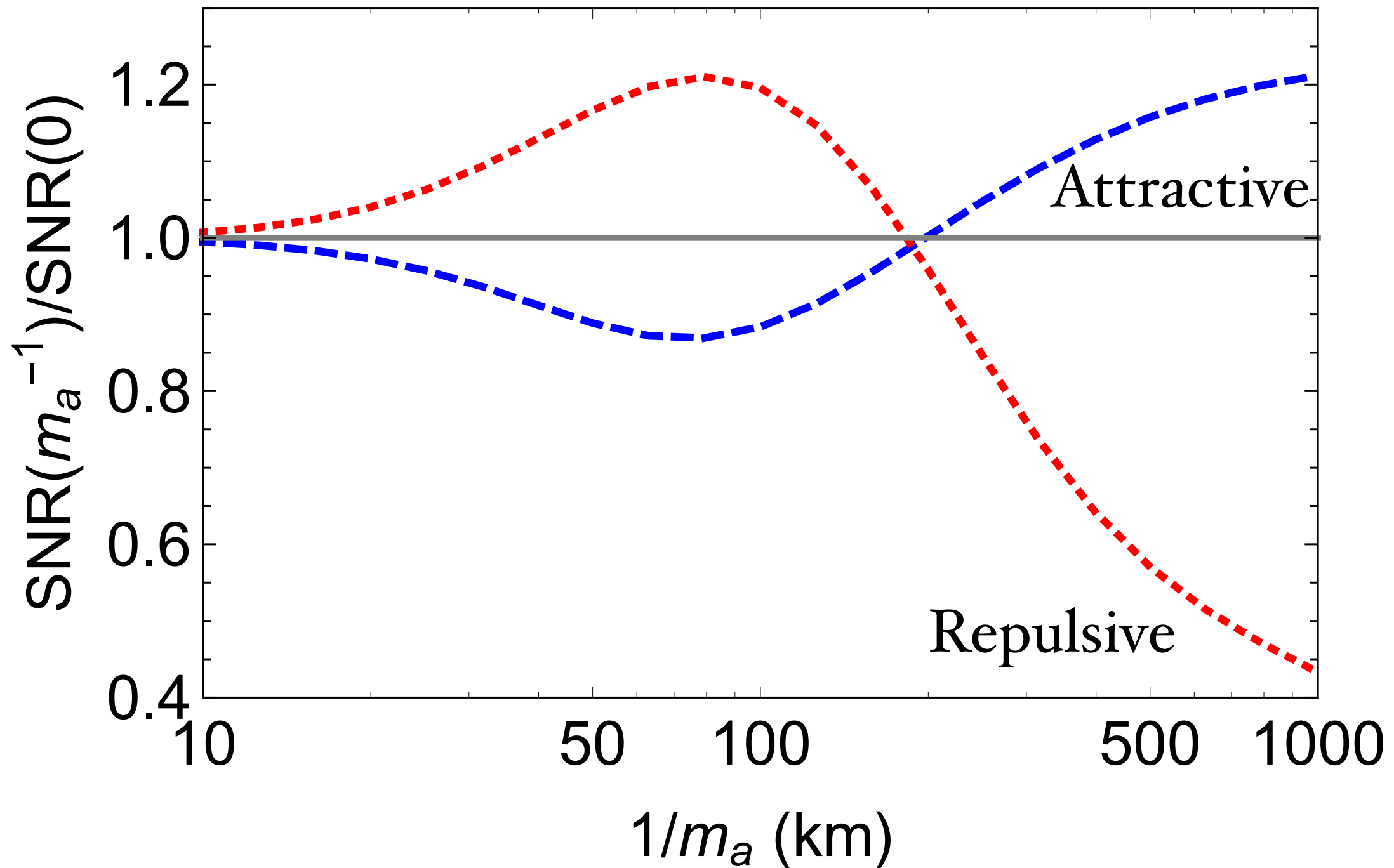
$$h_+(t) = \frac{4G\mu\omega^2 D^2}{r} \frac{1 + \cos \theta_i}{2} \cos 2\omega t,$$

$$h_\times(t) = \frac{4G\mu\omega^2 D^2}{r} \cos \theta_i \sin 2\omega t$$

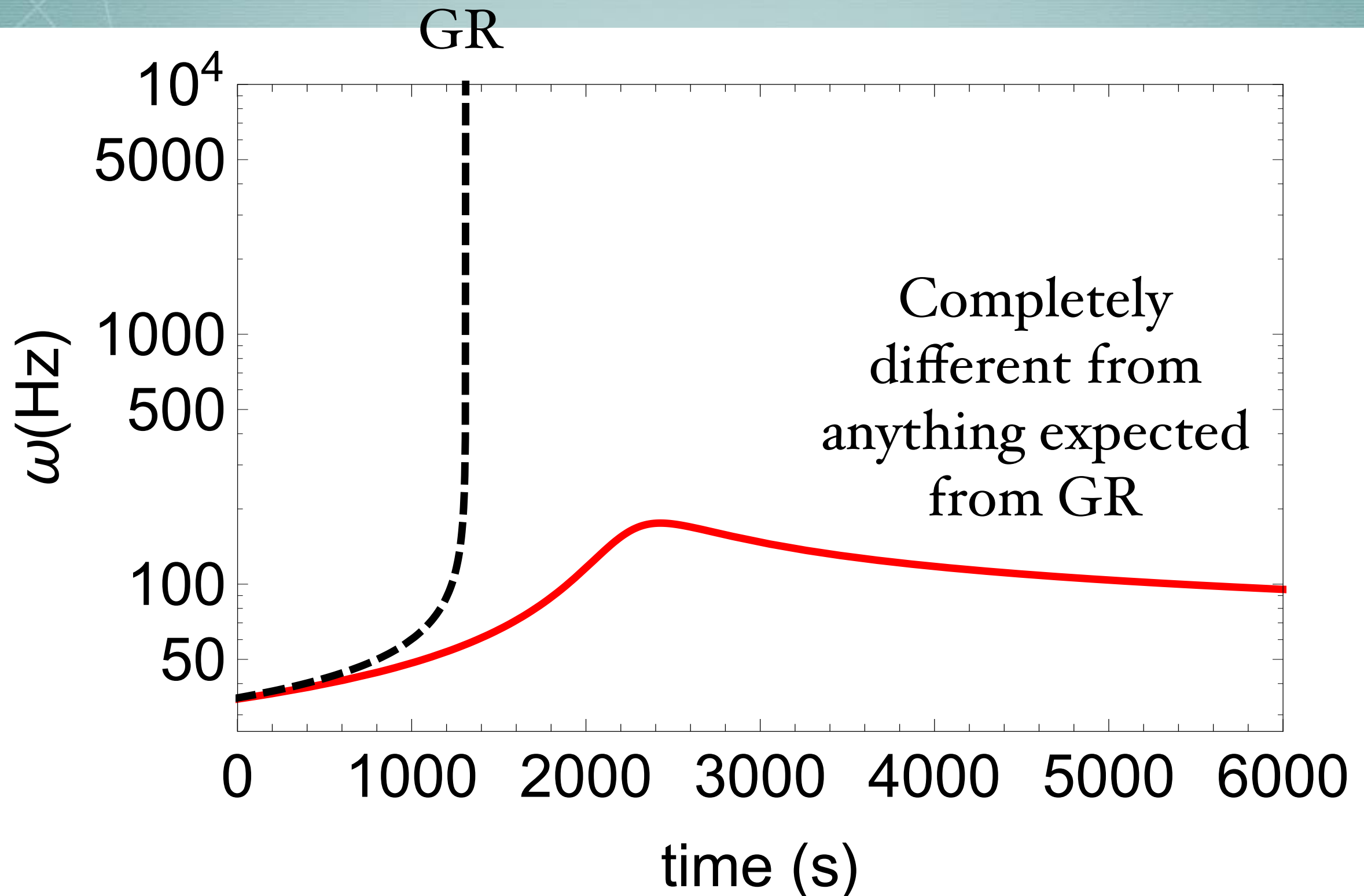
Inspiral - Optimal Filter



Inspiral - Optimal Filter



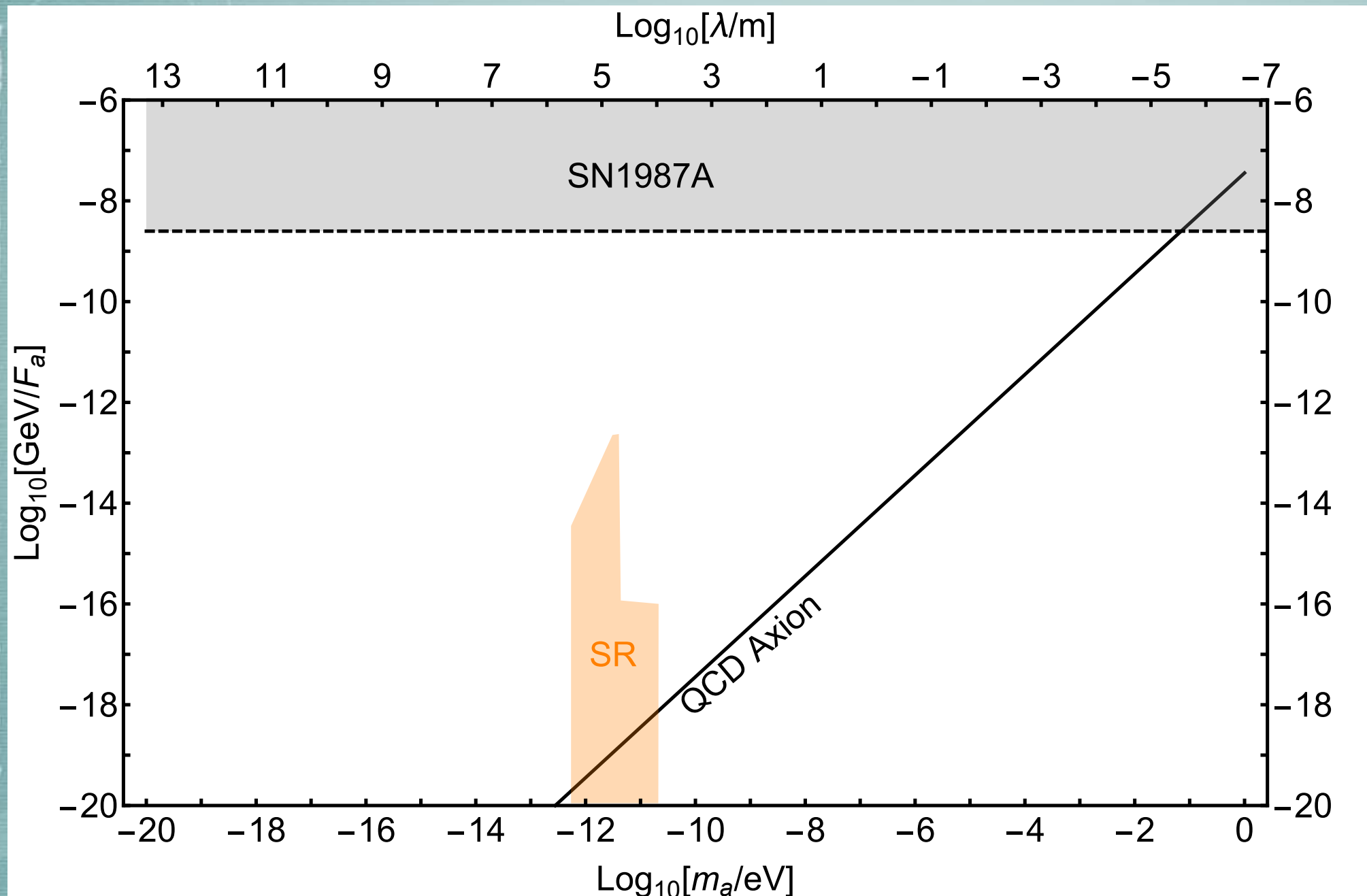
Inspiral - 110% gravity



Constraints

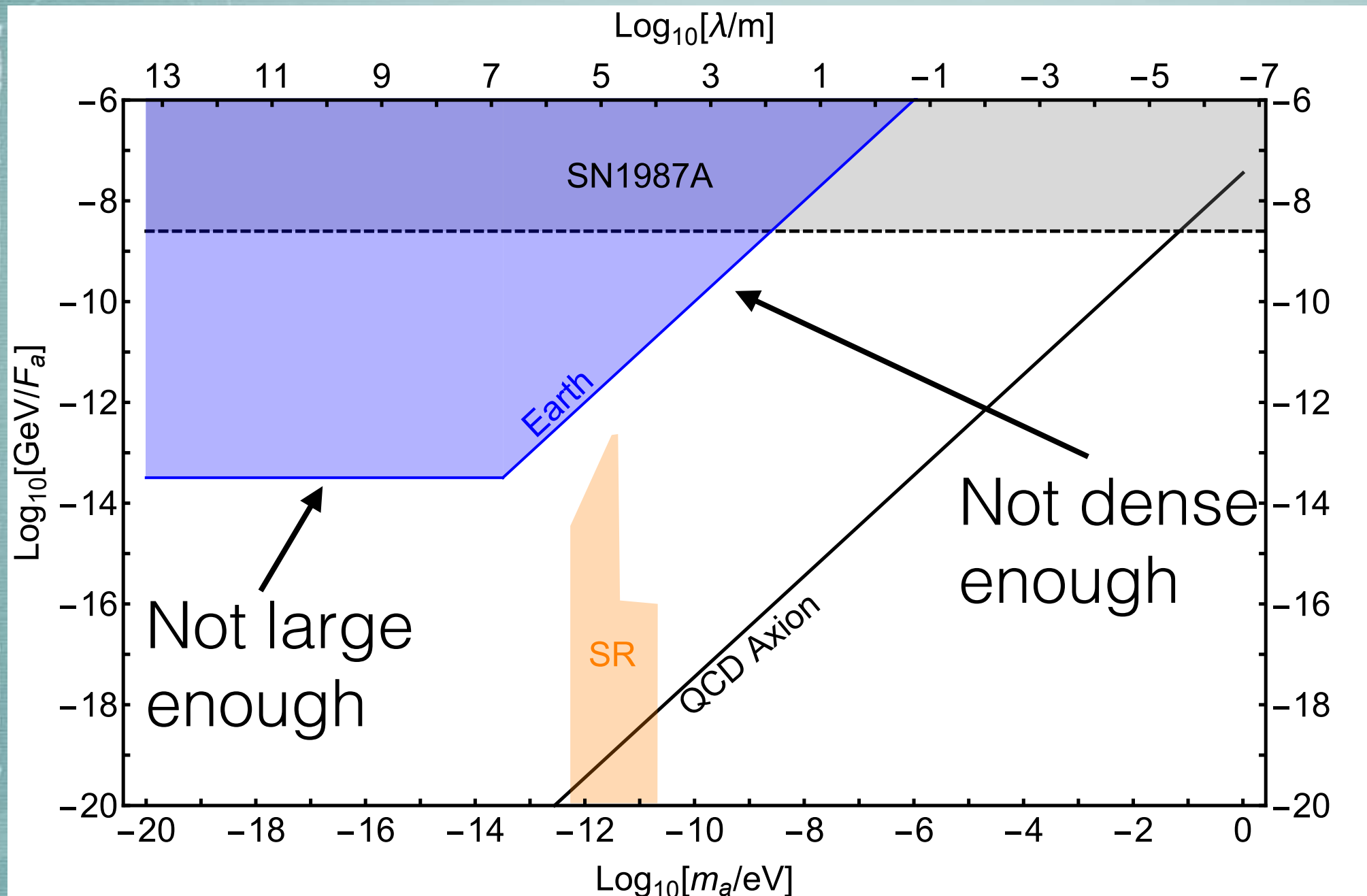
- As mentioned before objects can source the axion which leads to new constraints as theta angle is around π for these objects

Constraints



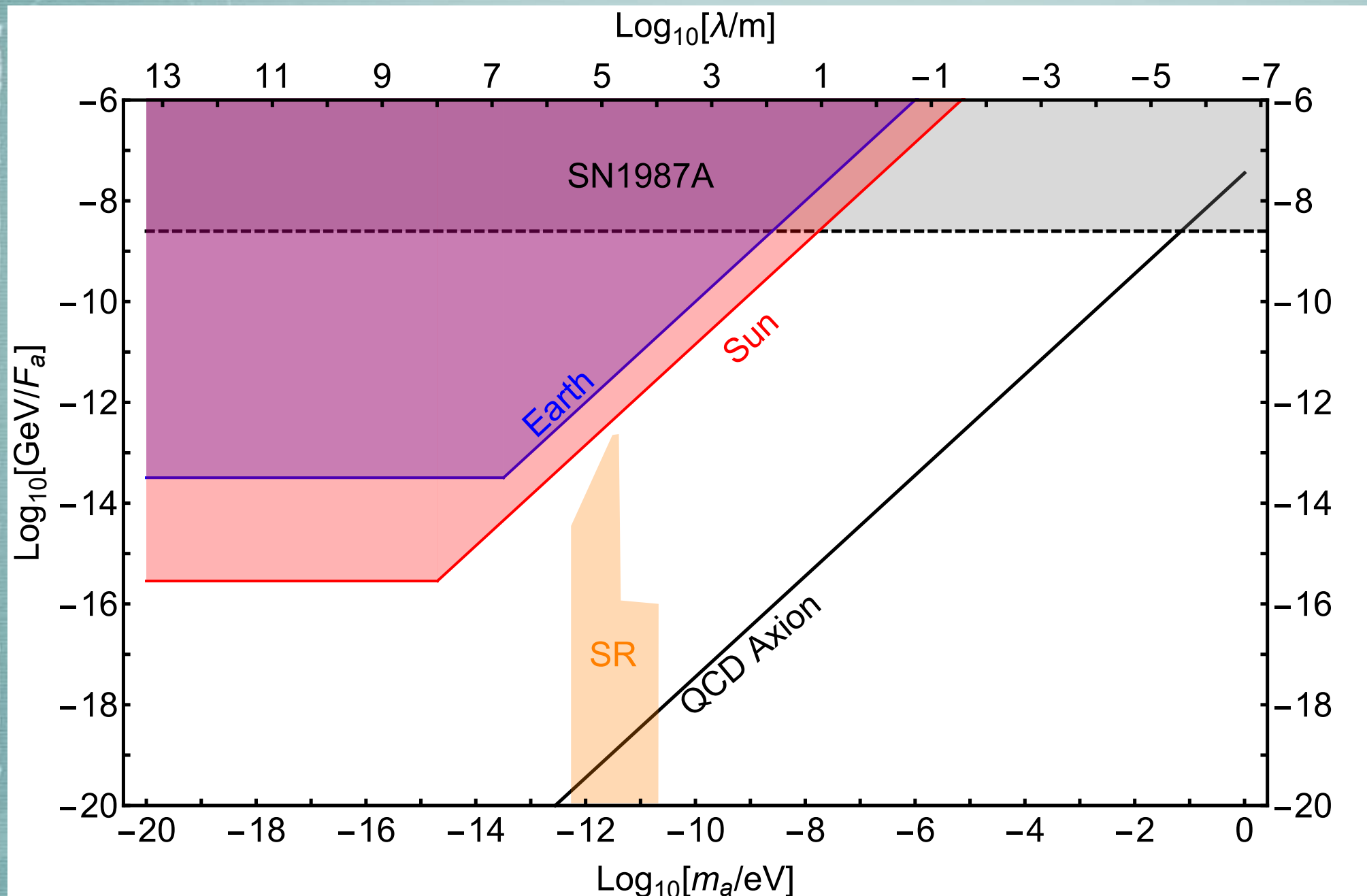
Super
Radiance
spin downs
highly
spinning
black holes

Constraints



Theta on
Earth
measured
to be very
small

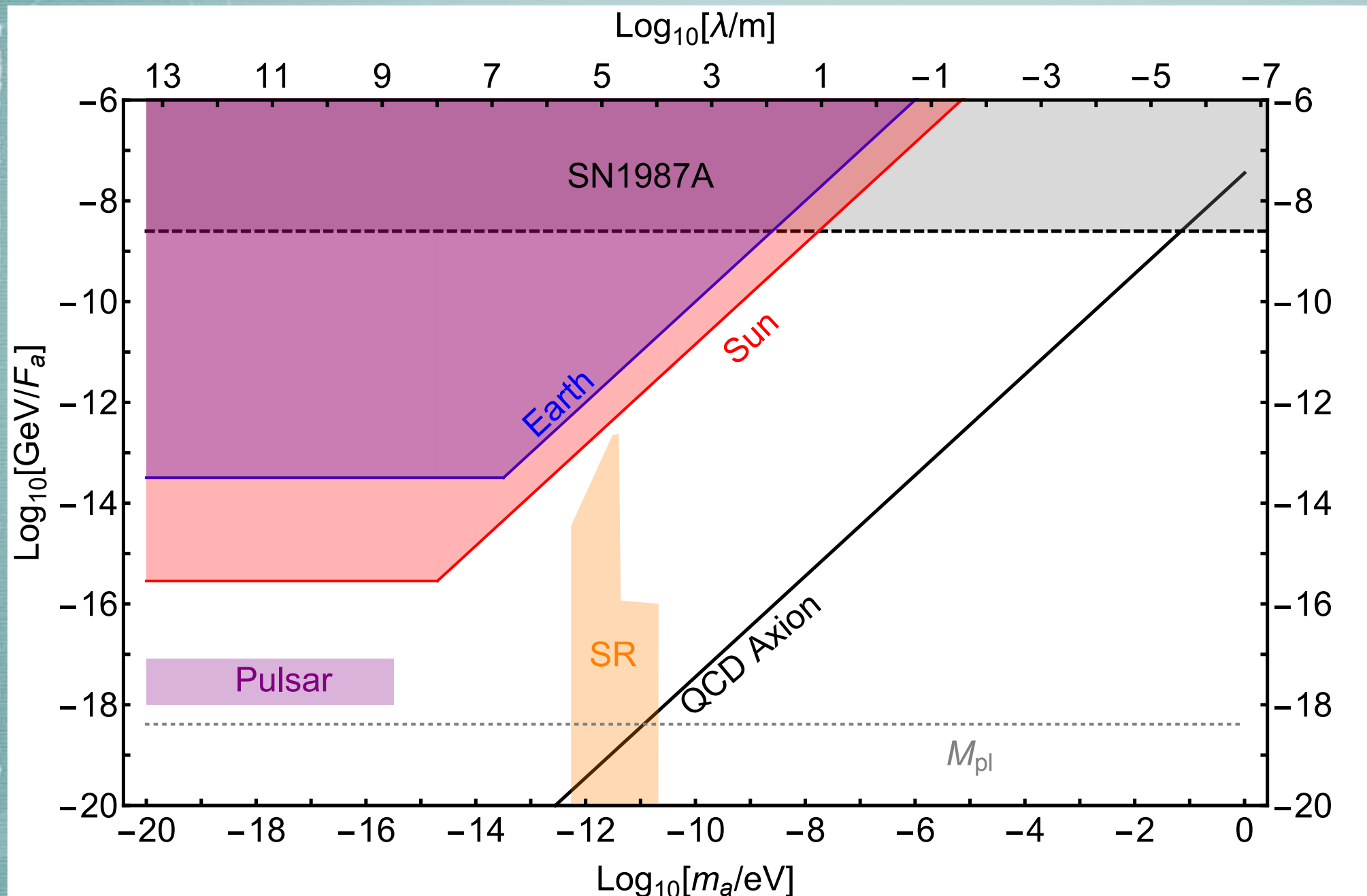
Constraints



$\text{Be}^7\text{-Li}^7$
862 keV
neutrino line
measured

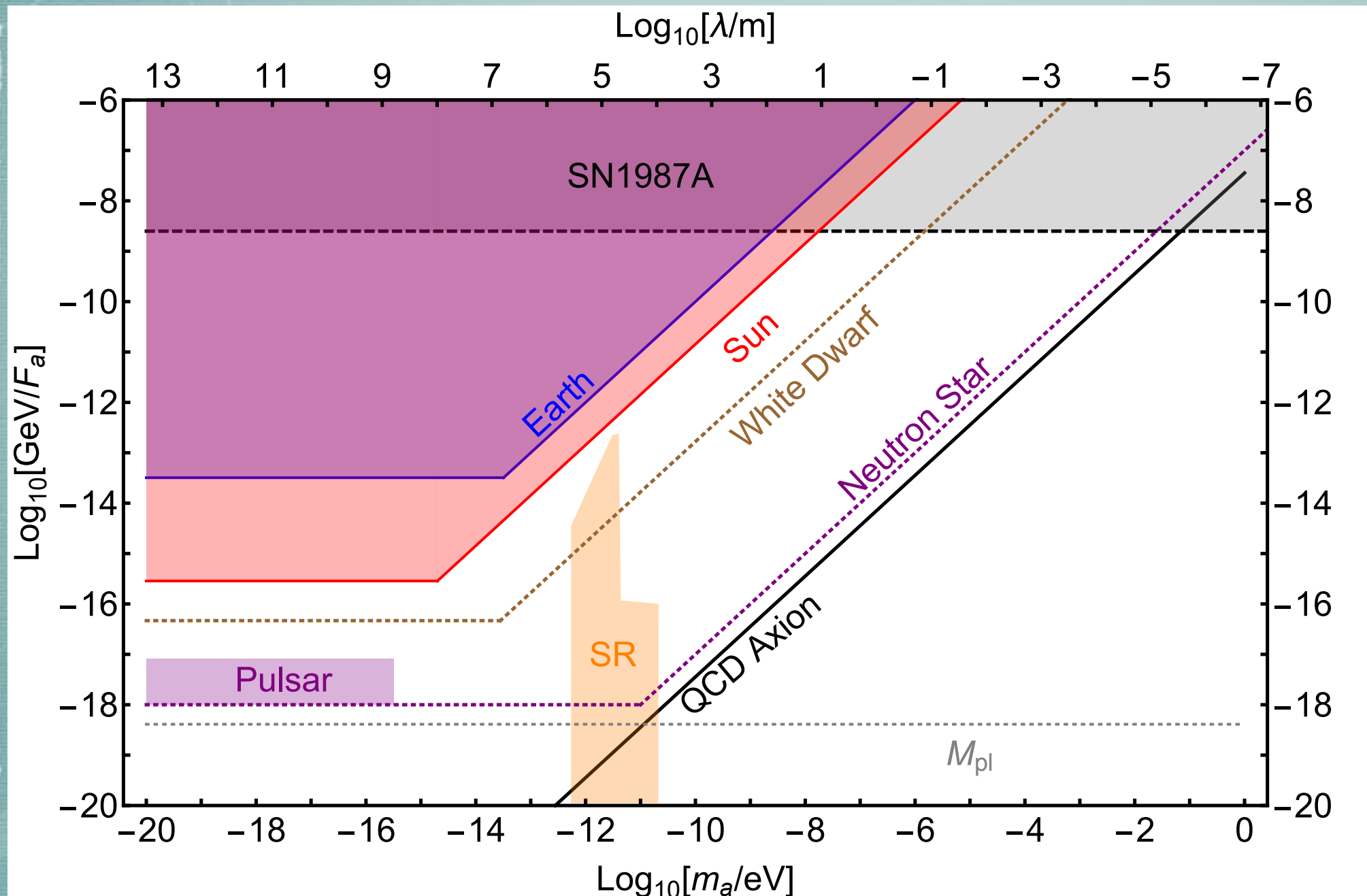
$m_p - m_n$
changed by
10 MeV

Constraints



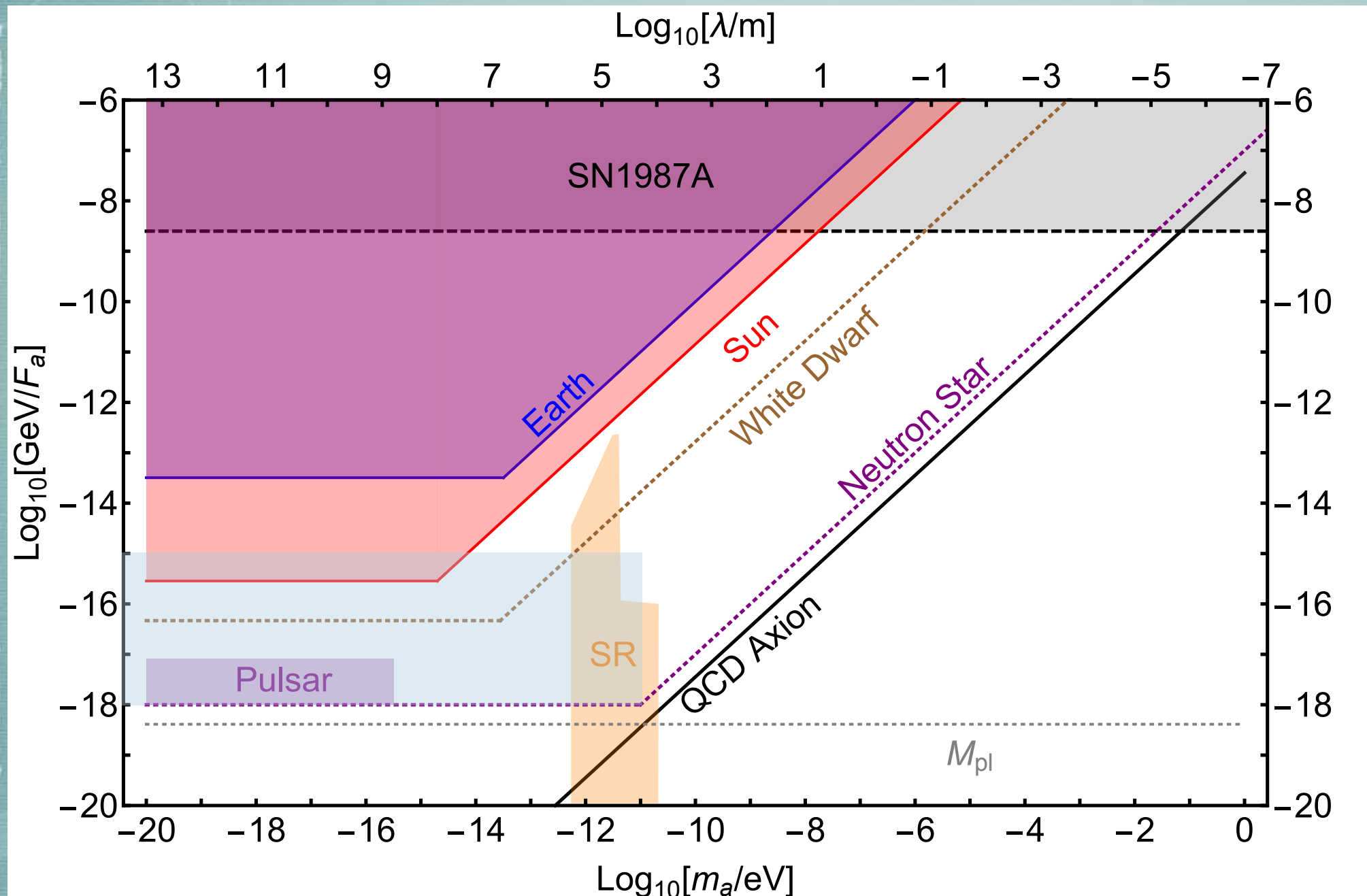
Hulse-Taylor
and other
binaries feel
only GR at
 $\mathcal{O}(1)$

Constraints



If we understood white dwarf/ neutron star nuclear emission lines at $\theta = \pi$

Constraints



Very dependent on equation of state and masses but in principle can exclude

Conclusion

- Tuning is not always bad
- Tuned axions can
 - Mediate forces stronger than gravity between neutron stars while evading 5th force experiments
 - Attractive or repulsive of force
 - Distance scale of force is naturally 10-100 km
 - Not necessarily a $1/r^2$ force

Conclusion

- To Do : Numerical Simulations!
 - To see that force is not $1/r^2$, neutron stars need to be very close
 - Interactions between neutron stars will change boundary condition
 - To see effect on late inspiral/merger
 - Effect of equation of state
 - What happens for a non tuned axion
 - ...